

# Measurement of the Splitting Function in $pp$ and Pb-Pb Collisions at $\sqrt{s_{NN}} = 5.02$ TeV

A. M. Sirunyan *et al.*\*  
(CMS Collaboration)

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Data from heavy ion collisions suggest that the evolution of a parton shower is modified by interactions with the color charges in the dense partonic medium created in these collisions, but it is not known where in the shower evolution the modifications occur. The momentum ratio of the two leading partons, resolved as subjects, provides information about the parton shower evolution. This substructure observable, known as the splitting function, reflects the process of a parton splitting into two other partons and has been measured for jets with transverse momentum between 140 and 500 GeV, in  $pp$  and PbPb collisions at a center-of-mass energy of 5.02 TeV per nucleon pair. In central PbPb collisions, the splitting function indicates a more unbalanced momentum ratio, compared to peripheral PbPb and  $pp$  collisions. The measurements are compared to various predictions from event generators and analytical calculations.

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Scattering processes with large momentum transfer  $Q$  between the partonic constituents of colliding nucleons occur early in heavy ion collisions. Further interactions of the outgoing partons with the produced (colored) hot and dense quantum chromodynamics (QCD) medium (the quark-gluon plasma, QGP) may modify the angular and momentum distributions of final-state hadronic jet fragments relative to those in proton-proton collisions. This process, known as jet quenching, can be used to probe the properties of the QGP [1,2]. Jet quenching was first observed at the Relativistic Heavy Ion Collider [3–9] and then at the Large Hadron Collider (LHC) [10–25]. This Letter reports an attempt to isolate parton splittings to two well separated partons with high transverse momentum ( $p_T$ ), probing medium induced effects during the parton shower evolution in the QGP. Information about these leading partons of a hard splitting can be obtained by removing the softer wide-angle radiation contributions, done through the use of jet grooming algorithms that attempt to split (“decluster”) a single jet into two subjects [26–30]. For a parton shower in vacuum, these subjects provide access to the properties of the first splitting in the parton evolution [31,32]. Interactions of the two outgoing partons with the QGP potentially modify the properties of subsequent splittings resulting in different subject properties. This Letter reports a study of hard parton splittings in  $pp$  and PbPb collisions.

An observable characterizing the parton splitting, denoted by  $z_g$ , is defined as the ratio between the  $p_T$  of the less energetic subject,  $p_{T,2}$ , and the  $p_T$  sum of the two subjects [32],  $z_g = p_{T,2}/(p_{T,1} + p_{T,2})$ . A measurement of the  $z_g$  distribution in  $pp$  collisions, using CMS open data, was recently reported [33,34]. In PbPb collisions, this measurement reflects how the two color-charged partons produced in the first splitting propagate through the QGP, probing the role of color coherence of the jet in the medium [35]. If the partons act as a single coherent emitter, the two subjects will be equally modified, leaving  $z_g$  unaffected [36]. If, instead, the partons in the medium act as decoherent emitters, the two subjects should be modified differently, thereby altering  $z_g$ . In addition,  $z_g$  is sensitive to semihard medium-induced gluon radiation [37], modifications of the initial parton splitting [38], and the medium response [39].

The analysis uses data collected by the CMS experiment in 2015. The PbPb and  $pp$  data samples, both at a nucleon-nucleon center-of-mass energy of 5.02 TeV, correspond to integrated luminosities of  $404 \mu\text{b}^{-1}$  and  $27.4 \text{ pb}^{-1}$ , respectively. The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter, and a brass and scintillator hadron calorimeter, each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity,  $\eta$ , coverage provided by the barrel and endcap detectors. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [40].

The particle-flow (PF) algorithm reconstructs and identifies each individual particle with an optimized combination of information from the various elements of the CMS detector [41]. The PF candidates identified as a photon or a

\*Full author list given at the end of the article.

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neutral hadron are treated as massless, while for charged hadrons the pion mass is assumed. The electron and muon PF candidates are assigned the corresponding lepton masses. Jets are reconstructed from the PF candidates using the anti- $k_T$  jet algorithm [42–44] with a distance parameter  $R = 0.4$ . The kinematics of the jet are determined using the vectorial sum of all particle momenta in the jet. For this analysis, jets are required to have  $p_{T,\text{jet}} > 140$  GeV and  $|\eta| < 1.3$ .

The online event selection trigger also uses the anti- $k_T$  algorithm with  $R = 0.4$  but applies a lower threshold on  $p_{T,\text{jet}}$ ; all events with a PF jet with  $p_{T,\text{jet}} > 80$  GeV were recorded in the  $pp$  case, while in PbPb collisions the triggers (based on jets reconstructed from calorimeter deposits including a subtraction for the uncorrelated underlying event) use a 100 GeV threshold. Noncollision events, such as beam-gas interactions or cosmic-ray muons, are rejected offline [19]. The events are required to have a primary vertex reconstructed within 15 cm (0.15 cm) of the nominal interaction point along the beam direction (in the transverse plane). The average number of additional collisions per bunch crossing is less than 0.9 in both data sets, having a negligible effect on the measurement. The PbPb event sample is divided into centrality intervals, reflecting the impact parameter of the colliding nuclei, using the percentage of the total inelastic hadronic cross section, which is evaluated using the sum of the total energy deposited in both forward hadron calorimeters, covering the  $3 < |\eta| < 5$  range [45].

The PYTHIA 6.423 [46] event generator (tune Z2\* [47,48]) is used to calculate Monte Carlo (MC) corrections. For PbPb simulations, the PYTHIA 6 events are embedded into an underlying event produced with HYDJET 1.9 [49]. All generated events undergo a full GEANT4 [50] simulation of the CMS detector response. Additional cross check samples are produced with PYTHIA 8.212 [51] (tune CUETP8M1 [48]) and HERWIG++ [52] (tune EE5C [53]).

In PbPb collisions, the constituents of the jet are corrected for the underlying event contribution using the “constituent subtraction” method [54], a particle-by-particle approach that removes or corrects jet constituents based on the average underlying event density. The subtraction corrects both the four-momentum of the jet and its substructure. Underlying event densities are determined by calculating the median  $p_T$  per unit area,  $\rho$ , and a density term related to the jet mass,  $\rho_m$ , using a procedure in which all of the particles in the event are clustered into jets using the  $k_T$  algorithm with  $R = 0.4$  [42,43,55]. To match the jets used in this analysis, only  $k_T$  jets with  $|\eta| < 1.3$  are included in the density determination. The influence of true hard jet fragments on the background estimation is reduced by excluding the two leading  $k_T$  jets. The constituent-subtracted jets are corrected for the detector response with jet energy corrections derived from independent  $pp$  and PbPb simulations. Additional corrections

for the mismodeling of the detector response are also applied [56].

Jet grooming algorithms aim to isolate the hard prongs of a jet and remove soft wide-angle radiation. The “soft drop” declustering procedure, used in this analysis, is an extension of the modified mass drop tagger [29]. The procedure starts by selecting an anti- $k_T$  jet that has already been constituent-subtracted and reclustered with the Cambridge-Aachen algorithm [57] to form a pairwise clustering tree with an angular-ordered structure. A pairwise declustering is performed on this tree. In each step of the declustering, a branching into two subjets is accepted if they pass the soft drop condition [30],

$$\frac{\min(p_{T,i}, p_{T,j})}{p_{T,i} + p_{T,j}} > z_{\text{cut}} \left( \frac{\Delta R_{ij}}{R_0} \right)^\beta, \quad (1)$$

where the subscripts “ $i$ ” and “ $j$ ” indicate the subjets at that step of the declustering,  $\Delta R_{ij}$  is the distance between the two subjets in the  $\eta$ - $\phi$  plane,  $R_0$  is the cone size of the anti- $k_T$  jet, and  $z_{\text{cut}}$  is an adjustable parameter. If the soft drop condition is not satisfied, the softer subjet is dropped. For this study,  $z_{\text{cut}}$  is set to 0.1 [30]. The parameter  $\beta$  is set to 0, which satisfies an extended version of infrared and collinear safety by absorbing the collinear divergences into a generalized fragmentation function recovering the QCD splitting function [32]. Once the soft drop condition is satisfied, the two subjets at that position in the tree are used in the analysis. If the soft drop condition is never satisfied, the jet is not used. This is the case for 1.5% of the jets measured at  $p_{T,\text{jet}} = 140$  GeV, increasing to 3.0% at  $p_{T,\text{jet}} = 300$  GeV, independent of collision centrality.

Groomed jets with a small distance between the two subjets frequently result from the ambiguous case where the two subjets cannot be distinctly resolved, leading to a significant misassignment of particle constituents to subjets. An additional selection of  $\Delta R_{12} > 0.1$  is applied, removing 40% (60%) of the jets measured at low (high)  $p_{T,\text{jet}}$ , to avoid an unphysical modification of  $z_g$ . This selection rejects an additional 15% (5%) of the jets at low (high)  $p_{T,\text{jet}}$  in the 10% most central PbPb collisions, in comparison to the noncentral collisions, an effect well reproduced by the simulation. The systematic uncertainty on the  $z_g$  variable is evaluated by varying the  $\Delta R_{12}$  minimum distance requirement by its one standard deviation MC resolution of 10%; this variation results in a 2% uncertainty, independent of centrality.

The transverse momentum of the jet after grooming,  $p_{T,g}$ , is identical to or smaller than the original  $p_{T,\text{jet}}$ . The groomed  $p_T$  fraction,  $p_{T,g}/p_{T,\text{jet}}$ , is compared to simulations in Fig. 1 for jets with  $160 < p_{T,\text{jet}} < 180$  GeV, in  $pp$  and central PbPb collisions. The measured and simulated distributions are in agreement.

The potential bias due to the online jet trigger is evaluated by using events collected with a lower threshold

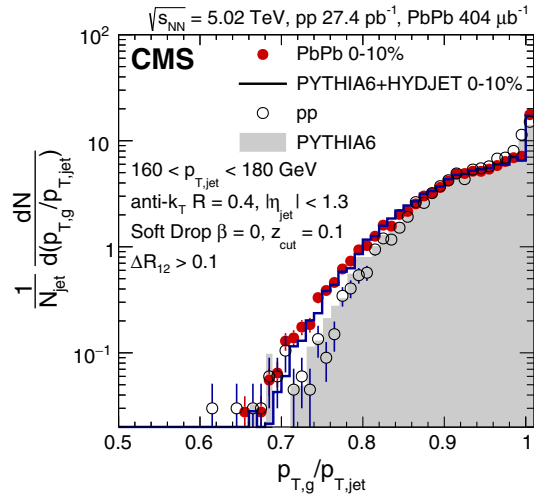


FIG. 1. Groomed jet energy fraction in  $pp$  and in the 10% most central PbPb collisions, for jets with  $160 < p_{T,\text{jet}} < 180$  GeV and  $|\eta_{\text{jet}}| < 1.3$ . The  $pp$  (PbPb) data are compared to PYTHIA 6 (embedded in HYDJET) distributions.

and also minimum bias events. For the 10% most central PbPb collisions, a bias is found in the lowest  $p_{T,\text{jet}}$  range,  $140 < p_{T,\text{jet}} < 160$  GeV, changing the yield by values linearly decreasing from +6% at  $z_g = 0.1$  to -15% at  $z_g = 0.5$ . In the 10%–30% centrality class, the bias is half as large, and it vanishes for more peripheral events. The full bias is corrected for and the magnitude of the correction is treated as a  $z_g$  systematic uncertainty. The trigger has no effect on the measurements at higher  $p_{T,\text{jet}}$ .

The systematic uncertainty in the jet energy scale, on the measured and simulated distributions, is obtained by propagating the uncertainties in the jet response correction [56,58]. A maximum deviation in yield of 4% is found in central PbPb collisions, decreasing to 2% in  $pp$  and peripheral PbPb collisions. This effect tends to increase (decrease) the  $p_T$  of the leading (subleading) subjet. The systematic uncertainty in the normalization of the  $z_g$  distributions is estimated to be 5% (3%) in central (peripheral) collisions. The relative uncertainty in the jet energy resolution is 10%, leading to an uncertainty smaller than 0.5% on the  $z_g$  distribution.

Figure 2 shows the  $z_g$  distribution measured in  $pp$  collisions, together with results obtained with PYTHIA 6, PYTHIA 8, and HERWIG++, including a full simulation of detector effects. Both PYTHIA simulations have a slightly steeper  $z_g$  distribution than the data, while HERWIG++ shows an opposite trend.

To compare the  $z_g$  distribution in  $pp$  and PbPb collisions, in given  $p_{T,\text{jet}}$  and centrality ranges, the measurements in  $pp$  collisions are adjusted to match the subjet resolution in PbPb data. The resolution correction is derived, for each  $p_{T,\text{jet}}$  and collision centrality range, from full detector simulation studies of the ratio of the  $z_g$  distributions between PYTHIA and PYTHIA embedded into

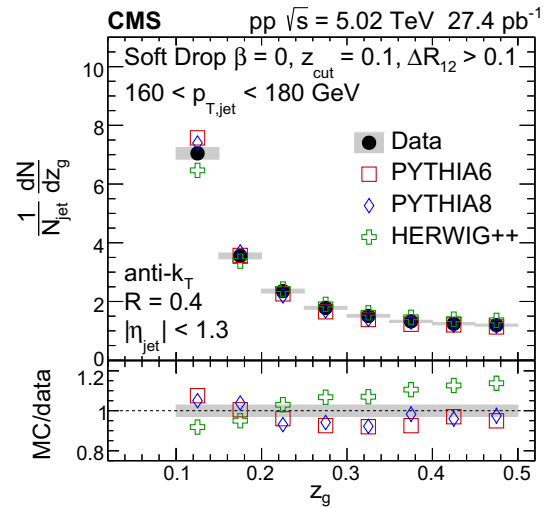


FIG. 2. The  $z_g$  distribution in  $pp$  collisions for  $160 < p_{T,\text{jet}} < 180$  GeV, compared to predictions from event generators. The error bars (shaded area) represent the statistical (systematic) uncertainty.

HYDJET. The ratio between simulated PbPb and  $pp$   $z_g$  distributions shows a relative decrease in the number of PbPb events at high  $z_g$ , reaching  $\sim 40\%$  in central collisions and negligible in peripheral collisions. The uncertainty in the correlation between the response of the two subjets is estimated by varying the individual subjet resolution by 10%, the relative correlation by 15%, and the subjet energy scale by 5%, corresponding to one standard deviation in resolution. This results in an uncertainty of 8%–10% in  $z_g$ . The mismodeling of the  $z_g$  distribution in PYTHIA, evaluated by reweighting to the  $z_g$  measurement in  $pp$  collisions, adds an uncertainty of 4%–5%. These uncertainties are assigned to the “smeared”  $pp$  data points. The resolution correction is validated with a parametric resolution model that uses the jet resolution and a sampled  $z_g$  in each  $p_{T,\text{jet}}$  range, and recreates the correction function for each centrality selection by sampling the individual subjet resolutions.

Figure 3 shows the  $z_g$  distributions measured in PbPb collisions, for several centrality intervals, in comparison to the smeared  $pp$  reference data. The systematic uncertainties on the  $z_g$  distributions are fully correlated from point to point, resulting in an anticorrelated uncertainty on the self-normalized distributions, and are uncorrelated between the  $pp$  and PbPb data sets. The  $z_g$  distribution in peripheral PbPb collisions agrees with the  $pp$  reference, while the more central collisions exhibit a steeper  $z_g$  distribution. Differences between the  $z_g$  of quark- and gluon-initiated jets are found to be a few percent [32], so that the observed modification cannot be attributed to the flavor composition within a fixed  $p_{T,\text{jet}}$  interval. The observation indicates that the splitting into two branches becomes increasingly more unbalanced as the PbPb collisions become more central.

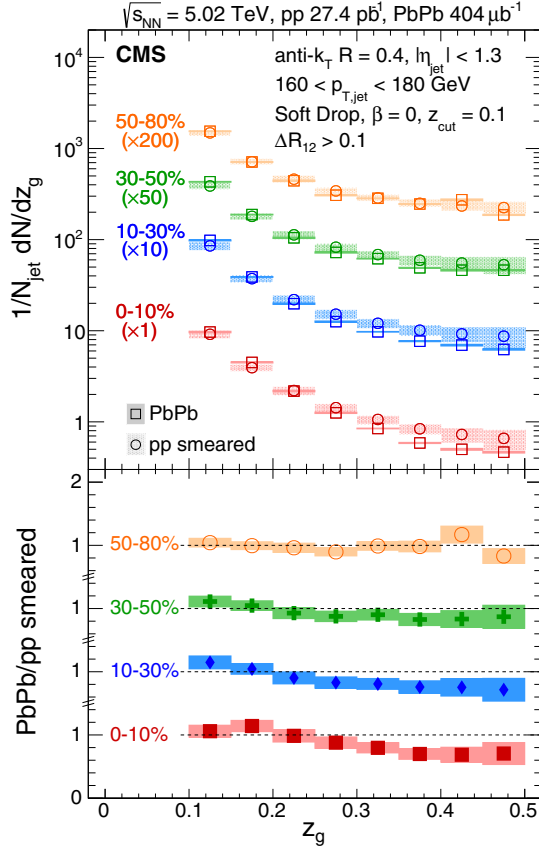


FIG. 3. The  $z_g$  distributions in PbPb collisions for  $160 < p_{T,\text{jet}} < 180$  GeV, in several centrality ranges, compared to  $pp$  data smeared to account for the differences in resolution. The error bars (shaded area) represent the statistical (systematic) uncertainty.

The modification of the  $z_g$  distribution in central PbPb collisions is shown in Fig. 4 over a wide kinematic range in  $p_{T,\text{jet}}$ . The measurement is compared to a prediction of the JEWEL event generator (shown with statistical and theoretical uncertainties originating from the treatment of the medium response), which incorporates medium-induced interactions while the partons propagate through the QGP [39,59,60]. The measurement is also compared with a soft-collinear effective theory (SCET) with Glauber gluon interactions [38] for two different quenching strengths, with a calculation incorporating multiple medium-induced gluon bremsstrahlung (BDMPS) [2,61,62] assuming that the two hard partons radiate gluons as a coherent emitter [37], and with a higher twist (HT) approach employing both coherent and incoherent energy loss [63]. Each of the three models is presented for two settings of the parameters reflecting their medium properties, as indicated in the legends, where  $L$  is the medium length,  $\hat{q}$  and  $\hat{q}_0$  denote medium transport coefficients, and  $g$  is the coupling strength between the jet and the medium. The BDMPS medium effect is too weak to describe the observed  $p_{T,\text{jet}}$  dependence, while the other models reproduce the data at low and high  $p_{T,\text{jet}}$ , using medium

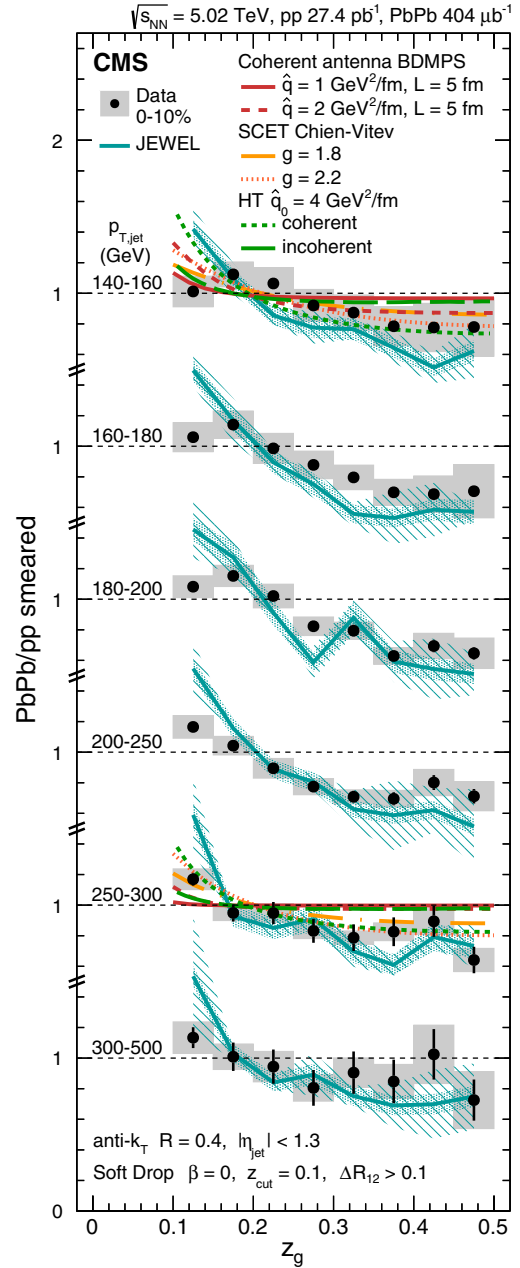


FIG. 4. Ratios of  $z_g$  distributions in PbPb and smeared  $pp$  collisions in the 10% most central events, for several  $p_{T,\text{jet}}$  ranges, compared to various jet quenching theoretical calculations [37–39,63]. The error bars (shaded area) represent the statistical (systematic) uncertainty. The diagonally hatched band denotes the uncertainty from the treatment of the medium response using the JEWEL event generator.

properties previously tuned to match measurements of the nuclear modification factors of charged hadrons and jets. For the HT calculation, the presence or absence of color coherence makes a significant difference. Since the detector resolution effects have a negligible impact on the theoretical calculations, given that they largely cancel in the PbPb to (smeared)  $pp$  ratio, the theoretical curves are shown without detector smearing.

In summary, the first measurement of the splitting function in  $pp$  and PbPb collisions at a center-of-mass energy of 5.02 TeV per nucleon pair has been presented. This represents the first application of a grooming technique to PbPb data, removing soft wide-angle radiation from the jet and thereby isolating the two leading subjects. The momentum sharing between these subjects is used to obtain information about hard parton splitting processes during the shower evolution. The PYTHIA and HERWIG++ event generators reproduce the measured splitting function in  $pp$  and peripheral PbPb collisions, at the level of 15%. In central PbPb collisions, a steeper  $z_g$  distribution is observed, indicating that the parton splitting process is modified by the hot medium created in heavy ion collisions. These results provide new insight into the role of color coherence and other attributes of the interactions of partons in the quark-gluon plasma.

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Janssen,<sup>4</sup> J. Lauwers,<sup>4</sup> H. Van Haevermaet,<sup>4</sup> P. Van Mechelen,<sup>4</sup> N. Van Remortel,<sup>4</sup> S. Abu Zeid,<sup>5</sup> F. Blekman,<sup>5</sup> J. D'Hondt,<sup>5</sup> I. De Bruyn,<sup>5</sup> J. De Clercq,<sup>5</sup> K. Deroover,<sup>5</sup> G. Flouris,<sup>5</sup> D. Lontkovskiy,<sup>5</sup> S. Lowette,<sup>5</sup> S. Moortgat,<sup>5</sup> L. Moreels,<sup>5</sup> Q. Python,<sup>5</sup> K. Skovpen,<sup>5</sup> S. Tavernier,<sup>5</sup> W. Van Doninck,<sup>5</sup> P. Van Mulders,<sup>5</sup> I. Van Parijs,<sup>5</sup> D. Beghin,<sup>6</sup> H. Brun,<sup>6</sup> B. Clerbaux,<sup>6</sup> G. De Lentdecker,<sup>6</sup> H. Delannoy,<sup>6</sup> B. Dorney,<sup>6</sup> G. Fasanella,<sup>6</sup> L. Favart,<sup>6</sup> R. Goldouzian,<sup>6</sup> A. Grebenyuk,<sup>6</sup> G. Karapostoli,<sup>6</sup> T. Lenzi,<sup>6</sup> J. Luetic,<sup>6</sup> T. Maerschalk,<sup>6</sup> A. Marinov,<sup>6</sup> A. Randle-conde,<sup>6</sup> T. Seva,<sup>6</sup> E. Starling,<sup>6</sup> C. Vander Velde,<sup>6</sup> P. Vanlaer,<sup>6</sup> D. Vannerom,<sup>6</sup> R. Yonamine,<sup>6</sup> F. Zenoni,<sup>6</sup> F. Zhang,<sup>6,c</sup> A. Cimmino,<sup>7</sup> T. Cornelis,<sup>7</sup> D. Dobur,<sup>7</sup> A. Fagot,<sup>7</sup> M. Gul,<sup>7</sup> I. Khvastunov,<sup>7,d</sup> D. Poyraz,<sup>7</sup> C. Roskas,<sup>7</sup> S. Salva,<sup>7</sup> M. Tytgat,<sup>7</sup> W. Verbeke,<sup>7</sup> N. Zaganidis,<sup>7</sup> H. Bakhshiansohi,<sup>8</sup> O. Bondu,<sup>8</sup> S. Brochet,<sup>8</sup> G. Bruno,<sup>8</sup> C. Caputo,<sup>8</sup> A. Caudron,<sup>8</sup> P. David,<sup>8</sup> S. De Visscher,<sup>8</sup> C. Delaere,<sup>8</sup> M. Delcourt,<sup>8</sup> B. Francois,<sup>8</sup> A. Giammanco,<sup>8</sup> M. Komm,<sup>8</sup> G. Krintiras,<sup>8</sup> V. Lemaître,<sup>8</sup> A. Magitteri,<sup>8</sup> A. Mertens,<sup>8</sup> M. Musich,<sup>8</sup> K. Piotrkowski,<sup>8</sup> L. Quertenmont,<sup>8</sup> A. Saggio,<sup>8</sup> M. Vidal Marono,<sup>8</sup> S. Wertz,<sup>8</sup> J. Zobec,<sup>8</sup> N. Belyi,<sup>9</sup> W. L. Aldá Júnior,<sup>10</sup> F. L. Alves,<sup>10</sup> G. A. Alves,<sup>10</sup> L. Brito,<sup>10</sup> M. Correa Martins Junior,<sup>10</sup> C. Hensel,<sup>10</sup> A. Moraes,<sup>10</sup> M. E. Pol,<sup>10</sup> P. Rebello Teles,<sup>10</sup> E. Belchior Batista Das Chagas,<sup>11</sup> W. Carvalho,<sup>11</sup> J. Chinellato,<sup>11,e</sup> E. Coelho,<sup>11</sup> E. M. Da Costa,<sup>11</sup> G. G. Da Silveira,<sup>11,f</sup> D. De Jesus Damiao,<sup>11</sup> S. Fonseca De Souza,<sup>11</sup> L. M. Huertas Guativa,<sup>11</sup> H. Malbouisson,<sup>11</sup> M. Melo De Almeida,<sup>11</sup> C. Mora Herrera,<sup>11</sup> L. Mundim,<sup>11</sup> H. Nogima,<sup>11</sup> L. J. Sanchez Rosas,<sup>11</sup> A. Santoro,<sup>11</sup> A. Sznajder,<sup>11</sup> M. Thiel,<sup>11</sup> E. J. Tonelli Manganote,<sup>11,e</sup> F. Torres Da Silva De Araujo,<sup>11</sup> A. Vilela Pereira,<sup>11</sup> S. Ahuja,<sup>12a</sup> C. A. Bernardes,<sup>12a</sup> T. R. Fernandez Perez Tomei,<sup>12a</sup> E. M. Gregores,<sup>12b</sup> P. G. Mercadante,<sup>12b</sup> S. F. Novaes,<sup>12a</sup> Sandra S. Padula,<sup>12a</sup> D. Romero Abad,<sup>12b</sup> J. C. Ruiz Vargas,<sup>12a</sup> A. Aleksandrov,<sup>13</sup> R. Hadjiiska,<sup>13</sup> P. Iaydjiev,<sup>13</sup> M. Misheva,<sup>13</sup> M. Rodozov,<sup>13</sup> M. Shopova,<sup>13</sup> G. Sultanov,<sup>13</sup> A. Dimitrov,<sup>14</sup> I. Glushkov,<sup>14</sup> L. Litov,<sup>14</sup> B. Pavlov,<sup>14</sup> P. Petkov,<sup>14</sup> W. Fang,<sup>15,g</sup> X. Gao,<sup>15,g</sup> L. Yuan,<sup>15</sup> M. Ahmad,<sup>16</sup> J. G. Bian,<sup>16</sup> G. M. Chen,<sup>16</sup> H. S. Chen,<sup>16</sup> M. Chen,<sup>16</sup> Y. Chen,<sup>16</sup> C. H. Jiang,<sup>16</sup> D. Leggat,<sup>16</sup> H. Liao,<sup>16</sup> Z. Liu,<sup>16</sup> F. Romeo,<sup>16</sup> S. M. Shaheen,<sup>16</sup> A. Spiezia,<sup>16</sup> J. Tao,<sup>16</sup> C. 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Ptochos,<sup>22</sup> P. A. Razis,<sup>22</sup> H. Rykaczewski,<sup>22</sup> M. Finger,<sup>23,i</sup> M. Finger Jr.,<sup>23,i</sup> E. Carrera Jarrin,<sup>24</sup> A. A. Abdelalim,<sup>25,j,k</sup> Y. Mohammed,<sup>25,l</sup> E. Salama,<sup>25,m,n</sup> R. K. Dewanjee,<sup>26</sup> M. Kadastik,<sup>26</sup> L. Perrini,<sup>26</sup> M. Raidal,<sup>26</sup> A. Tiko,<sup>26</sup> C. Veelken,<sup>26</sup> P. Eerola,<sup>27</sup> H. Kirschenmann,<sup>27</sup> J. Pekkanen,<sup>27</sup> M. Voutilainen,<sup>27</sup> T. Järvinen,<sup>28</sup> V. Karimäki,<sup>28</sup> R. Kinnunen,<sup>28</sup> T. Lampén,<sup>28</sup> K. Lassila-Perini,<sup>28</sup> S. Lehti,<sup>28</sup> T. Lindén,<sup>28</sup> P. Luukka,<sup>28</sup> E. Tuominen,<sup>28</sup> J. Tuominiemi,<sup>28</sup> J. Talvitie,<sup>29</sup> T. Tuuva,<sup>29</sup> M. Besancon,<sup>30</sup> F. Couderc,<sup>30</sup> M. Dejardin,<sup>30</sup> D. Denegri,<sup>30</sup> J. L. Faure,<sup>30</sup> F. Ferri,<sup>30</sup> S. Ganjour,<sup>30</sup> S. Ghosh,<sup>30</sup> A. Givernaud,<sup>30</sup> P. Gras,<sup>30</sup> G. Hamel de Monchenault,<sup>30</sup> P. Jarry,<sup>30</sup> I. Kucher,<sup>30</sup> C. Leloup,<sup>30</sup> E. Locci,<sup>30</sup> M. Machet,<sup>30</sup> J. Malcles,<sup>30</sup> G. Negro,<sup>30</sup> J. Rander,<sup>30</sup> A. Rosowsky,<sup>30</sup> M. Ö. Sahin,<sup>30</sup> M. Titov,<sup>30</sup> A. Abdulsalam,<sup>31</sup> C. Amendola,<sup>31</sup> I. Antropov,<sup>31</sup> S. Baffioni,<sup>31</sup> F. Beaudette,<sup>31</sup> P. Busson,<sup>31</sup> L. Cadamuro,<sup>31</sup> C. Charlot,<sup>31</sup> R. Granier de Cassagnac,<sup>31</sup> M. Jo,<sup>31</sup> S. Lisniak,<sup>31</sup> A. Lobanov,<sup>31</sup> J. Martin Blanco,<sup>31</sup> M. Nguyen,<sup>31</sup> C. Ochando,<sup>31</sup> G. Ortona,<sup>31</sup> P. Paganini,<sup>31</sup> P. Pigard,<sup>31</sup> R. Salerno,<sup>31</sup> J. B. Sauvan,<sup>31</sup> Y. Sirois,<sup>31</sup> A. G. Stahl Leiton,<sup>31</sup> T. Strebler,<sup>31</sup> Y. Yilmaz,<sup>31</sup> A. Zabi,<sup>31</sup> A. Zghiche,<sup>31</sup> J.-L. Agram,<sup>32,o</sup> J. Andrea,<sup>32</sup> D. Bloch,<sup>32</sup> J.-M. Brom,<sup>32</sup> M. Buttignol,<sup>32</sup> E. C. Chabert,<sup>32</sup> N. Chanon,<sup>32</sup> C. Collard,<sup>32</sup> E. Conte,<sup>32,o</sup> X. Coubez,<sup>32</sup> J.-C. Fontaine,<sup>32,o</sup> D. Gelé,<sup>32</sup> U. Goerlach,<sup>32</sup> M. Jansová,<sup>32</sup> A.-C. Le Bihan,<sup>32</sup> N. Tonon,<sup>32</sup> P. Van Hove,<sup>32</sup> S. Gadrat,<sup>33</sup> S. Beauceron,<sup>34</sup> C. Bernet,<sup>34</sup> G. Boudoul,<sup>34</sup> R. Chierici,<sup>34</sup> D. Contardo,<sup>34</sup> P. Depasse,<sup>34</sup> H. El Mamouni,<sup>34</sup> J. Fay,<sup>34</sup> L. Finco,<sup>34</sup> S. Gascon,<sup>34</sup> M. Gouzevitch,<sup>34</sup>

G. Grenier,<sup>34</sup> B. Ille,<sup>34</sup> F. Lagarde,<sup>34</sup> I. B. Laktineh,<sup>34</sup> M. Lethuillier,<sup>34</sup> L. Mirabito,<sup>34</sup> A. L. Pequegnot,<sup>34</sup> S. Perries,<sup>34</sup>  
 A. Popov,<sup>34,p</sup> V. Sordini,<sup>34</sup> M. Vander Donckt,<sup>34</sup> S. Viret,<sup>34</sup> A. Khvedelidze,<sup>35,i</sup> Z. Tsamalaidze,<sup>36,i</sup> C. Autermann,<sup>37</sup>  
 L. Feld,<sup>37</sup> M. K. Kiesel,<sup>37</sup> K. Klein,<sup>37</sup> M. Lipinski,<sup>37</sup> M. Preuten,<sup>37</sup> C. Schomakers,<sup>37</sup> J. Schulz,<sup>37</sup> V. Zhukov,<sup>37,p</sup> A. Albert,<sup>38</sup>  
 E. Dietz-Laursonn,<sup>38</sup> D. Duchardt,<sup>38</sup> M. Endres,<sup>38</sup> M. Erdmann,<sup>38</sup> S. Erdweg,<sup>38</sup> T. Esch,<sup>38</sup> R. Fischer,<sup>38</sup> A. Güth,<sup>38</sup>  
 M. Hamer,<sup>38</sup> T. Hebbeker,<sup>38</sup> C. Heidemann,<sup>38</sup> K. Hoepfner,<sup>38</sup> S. Knutzen,<sup>38</sup> M. Merschmeyer,<sup>38</sup> A. Meyer,<sup>38</sup> P. Millet,<sup>38</sup>  
 S. Mukherjee,<sup>38</sup> T. Pook,<sup>38</sup> M. Radziej,<sup>38</sup> H. Reithler,<sup>38</sup> M. Rieger,<sup>38</sup> F. Scheuch,<sup>38</sup> D. Teyssier,<sup>38</sup> S. Thüer,<sup>38</sup> G. Flügge,<sup>39</sup>  
 B. Kargoll,<sup>39</sup> T. Kress,<sup>39</sup> A. Künsken,<sup>39</sup> T. Müller,<sup>39</sup> A. Nehr Korn,<sup>39</sup> A. Nowack,<sup>39</sup> C. Pistone,<sup>39</sup> O. Pooth,<sup>39</sup> A. Stahl,<sup>39,q</sup>  
 M. Aldaya Martin,<sup>40</sup> T. Arndt,<sup>40</sup> C. Asawatangtrakuldee,<sup>40</sup> K. Beernaert,<sup>40</sup> O. Behnke,<sup>40</sup> U. Behrens,<sup>40</sup>  
 A. Bermúdez Martínez,<sup>40</sup> A. A. Bin Anuar,<sup>40</sup> K. Borras,<sup>40,r</sup> V. Botta,<sup>40</sup> A. Campbell,<sup>40</sup> P. Connor,<sup>40</sup> C. Contreras-Campana,<sup>40</sup>  
 F. Costanza,<sup>40</sup> C. Diez Pardos,<sup>40</sup> G. Eckerlin,<sup>40</sup> D. Eckstein,<sup>40</sup> T. Eichhorn,<sup>40</sup> E. Eren,<sup>40</sup> E. Gallo,<sup>40,s</sup> J. Garay Garcia,<sup>40</sup>  
 A. Geiser,<sup>40</sup> A. Gizhko,<sup>40</sup> J. M. Grados Luyando,<sup>40</sup> A. Grohsjean,<sup>40</sup> P. Gunnellini,<sup>40</sup> M. Guthoff,<sup>40</sup> A. Harb,<sup>40</sup> J. Hauk,<sup>40</sup>  
 M. Hempel,<sup>40,t</sup> H. Jung,<sup>40</sup> A. Kalogeropoulos,<sup>40</sup> M. Kasemann,<sup>40</sup> J. Keaveney,<sup>40</sup> C. Kleinwort,<sup>40</sup> I. Korol,<sup>40</sup> D. Krücker,<sup>40</sup>  
 W. Lange,<sup>40</sup> A. Lelek,<sup>40</sup> T. Lenz,<sup>40</sup> J. Leonard,<sup>40</sup> K. Lipka,<sup>40</sup> W. Lohmann,<sup>40,t</sup> R. Mankel,<sup>40</sup> I.-A. Melzer-Pellmann,<sup>40</sup>  
 A. B. Meyer,<sup>40</sup> G. Mittag,<sup>40</sup> J. Mnich,<sup>40</sup> A. Mussgiller,<sup>40</sup> E. Ntomari,<sup>40</sup> D. Pitzl,<sup>40</sup> A. Raspereza,<sup>40</sup> B. Roland,<sup>40</sup>  
 M. Savitskiy,<sup>40</sup> P. Saxena,<sup>40</sup> R. Shevchenko,<sup>40</sup> S. Spannagel,<sup>40</sup> N. Stefaniuk,<sup>40</sup> G. P. Van Onsem,<sup>40</sup> R. Walsh,<sup>40</sup> Y. Wen,<sup>40</sup>  
 K. Wichmann,<sup>40</sup> C. Wissing,<sup>40</sup> O. Zenaiev,<sup>40</sup> R. Aggleton,<sup>41</sup> S. Bein,<sup>41</sup> V. Blobel,<sup>41</sup> M. Centis Vignali,<sup>41</sup> T. Dreyer,<sup>41</sup>  
 E. Garutti,<sup>41</sup> D. Gonzalez,<sup>41</sup> J. Haller,<sup>41</sup> A. Hinzmann,<sup>41</sup> M. Hoffmann,<sup>41</sup> A. Karavdina,<sup>41</sup> R. Klanner,<sup>41</sup> R. Kogler,<sup>41</sup>  
 N. Kovalchuk,<sup>41</sup> S. Kurz,<sup>41</sup> T. Lapsien,<sup>41</sup> I. Marchesini,<sup>41</sup> D. Marconi,<sup>41</sup> M. Meyer,<sup>41</sup> M. Niedziela,<sup>41</sup> D. Nowatschin,<sup>41</sup>  
 F. Pantaleo,<sup>41,q</sup> T. Peiffer,<sup>41</sup> A. Perieanu,<sup>41</sup> C. Scharf,<sup>41</sup> P. Schleper,<sup>41</sup> A. Schmidt,<sup>41</sup> S. Schumann,<sup>41</sup> J. Schwandt,<sup>41</sup>  
 J. Sonneveld,<sup>41</sup> H. Stadie,<sup>41</sup> G. Steinbrück,<sup>41</sup> F. M. Stober,<sup>41</sup> M. Stöver,<sup>41</sup> H. Tholen,<sup>41</sup> D. Troendle,<sup>41</sup> E. Usai,<sup>41</sup>  
 L. Vanelderen,<sup>41</sup> A. Vanhoefer,<sup>41</sup> B. Vormwald,<sup>41</sup> M. Akbiyik,<sup>42</sup> C. Barth,<sup>42</sup> S. Baur,<sup>42</sup> E. Butz,<sup>42</sup> R. Caspart,<sup>42</sup> T. Chwalek,<sup>42</sup>  
 F. Colombo,<sup>42</sup> W. De Boer,<sup>42</sup> A. Dierlamm,<sup>42</sup> B. Freund,<sup>42</sup> R. Friese,<sup>42</sup> M. Giffels,<sup>42</sup> D. Haitz,<sup>42</sup> M. A. Harrendorf,<sup>42</sup>  
 F. Hartmann,<sup>42,q</sup> S. M. Heindl,<sup>42</sup> U. Husemann,<sup>42</sup> F. Kassel,<sup>42,q</sup> S. Kudella,<sup>42</sup> H. Mildner,<sup>42</sup> M. U. Mozer,<sup>42</sup> Th. Müller,<sup>42</sup>  
 M. Plagge,<sup>42</sup> G. Quast,<sup>42</sup> K. Rabbertz,<sup>42</sup> M. Schröder,<sup>42</sup> I. Shvetsov,<sup>42</sup> G. Sieber,<sup>42</sup> H. J. Simonis,<sup>42</sup> R. Ulrich,<sup>42</sup> S. Wayand,<sup>42</sup>  
 M. Weber,<sup>42</sup> T. Weiler,<sup>42</sup> S. Williamson,<sup>42</sup> C. Wöhrmann,<sup>42</sup> R. Wolf,<sup>42</sup> G. Anagnostou,<sup>43</sup> G. Daskalakis,<sup>43</sup> T. Geralis,<sup>43</sup>  
 V. A. Giakoumopoulou,<sup>43</sup> A. Kyriakis,<sup>43</sup> D. Loukas,<sup>43</sup> I. Topsis-Giotis,<sup>43</sup> G. Karathanasis,<sup>44</sup> S. Kesisoglou,<sup>44</sup>  
 A. Panagiotou,<sup>44</sup> N. Saoulidou,<sup>44</sup> K. Kousouris,<sup>45</sup> I. Evangelou,<sup>46</sup> C. Foudas,<sup>46</sup> P. Kokkas,<sup>46</sup> S. Mallios,<sup>46</sup> N. Manthos,<sup>46</sup>  
 I. Papadopoulos,<sup>46</sup> E. Paradas,<sup>46</sup> J. Strologas,<sup>46</sup> F. A. Triantis,<sup>46</sup> M. Csanad,<sup>47</sup> N. Filipovic,<sup>47</sup> G. Pasztor,<sup>47</sup> O. Surányi,<sup>47</sup>  
 G. I. Veres,<sup>47,u</sup> G. Bencze,<sup>48</sup> C. Hajdu,<sup>48</sup> D. Horvath,<sup>48,v</sup> Á. Hunyadi,<sup>48</sup> F. Sikler,<sup>48</sup> V. Veszpremi,<sup>48</sup> A. J. Zsigmond,<sup>48</sup>  
 N. Beni,<sup>49</sup> S. Czellar,<sup>49</sup> J. Karancsi,<sup>49,w</sup> A. Makovec,<sup>49</sup> J. Molnar,<sup>49</sup> Z. Szillasi,<sup>49</sup> M. Bartók,<sup>50,u</sup> P. Raics,<sup>50</sup> Z. L. Trocsanyi,<sup>50</sup>  
 B. Ujvari,<sup>50</sup> S. Choudhury,<sup>51</sup> J. R. Komaragiri,<sup>51</sup> S. Bahinipati,<sup>52,x</sup> S. Bhowmik,<sup>52</sup> P. Mal,<sup>52</sup> K. Mandal,<sup>52</sup> A. Nayak,<sup>52,y</sup>  
 D. K. Sahoo,<sup>52,x</sup> N. Sahoo,<sup>52</sup> S. K. Swain,<sup>52</sup> S. Bansal,<sup>53</sup> S. B. Beri,<sup>53</sup> V. Bhatnagar,<sup>53</sup> R. Chawla,<sup>53</sup> N. Dhingra,<sup>53</sup>  
 A. K. Kalsi,<sup>53</sup> A. Kaur,<sup>53</sup> M. Kaur,<sup>53</sup> S. Kaur,<sup>53</sup> R. Kumar,<sup>53</sup> P. Kumari,<sup>53</sup> A. Mehta,<sup>53</sup> J. B. Singh,<sup>53</sup> G. Walia,<sup>53</sup>  
 Ashok Kumar,<sup>54</sup> Aashaq Shah,<sup>54</sup> A. Bhardwaj,<sup>54</sup> S. Chauhan,<sup>54</sup> B. C. Choudhary,<sup>54</sup> R. B. Garg,<sup>54</sup> S. Keshri,<sup>54</sup> A. Kumar,<sup>54</sup>  
 S. Malhotra,<sup>54</sup> M. Naimuddin,<sup>54</sup> K. Ranjan,<sup>54</sup> R. Sharma,<sup>54</sup> R. Bhardwaj,<sup>55</sup> R. Bhattacharya,<sup>55</sup> S. Bhattacharya,<sup>55</sup>  
 U. Bhawandeep,<sup>55</sup> S. Dey,<sup>55</sup> S. Dutt,<sup>55</sup> S. Dutta,<sup>55</sup> S. Ghosh,<sup>55</sup> N. Majumdar,<sup>55</sup> A. Modak,<sup>55</sup> K. Mondal,<sup>55</sup>  
 S. Mukhopadhyay,<sup>55</sup> S. Nandan,<sup>55</sup> A. Purohit,<sup>55</sup> A. Roy,<sup>55</sup> D. Roy,<sup>55</sup> S. Roy Chowdhury,<sup>55</sup> S. Sarkar,<sup>55</sup> M. Sharan,<sup>55</sup>  
 S. Thakur,<sup>55</sup> P. K. Behera,<sup>56</sup> R. Chudasama,<sup>57</sup> D. Dutta,<sup>57</sup> V. Jha,<sup>57</sup> V. Kumar,<sup>57</sup> A. K. Mohanty,<sup>57,q</sup> P. K. Netrakanti,<sup>57</sup>  
 L. M. Pant,<sup>57</sup> P. Shukla,<sup>57</sup> A. Topkar,<sup>57</sup> T. Aziz,<sup>58</sup> S. Dugad,<sup>58</sup> B. Mahakud,<sup>58</sup> S. Mitra,<sup>58</sup> G. B. Mohanty,<sup>58</sup> N. Sur,<sup>58</sup>  
 B. Sutar,<sup>58</sup> S. Banerjee,<sup>59</sup> S. Bhattacharya,<sup>59</sup> S. Chatterjee,<sup>59</sup> P. Das,<sup>59</sup> M. Guchait,<sup>59</sup> Sa. Jain,<sup>59</sup> S. Kumar,<sup>59</sup> M. Maity,<sup>59,z</sup>  
 G. Majumder,<sup>59</sup> K. Mazumdar,<sup>59</sup> T. Sarkar,<sup>59,z</sup> N. Wickramage,<sup>59,aa</sup> S. Chauhan,<sup>60</sup> S. Dube,<sup>60</sup> V. Hegde,<sup>60</sup> A. Kapoor,<sup>60</sup>  
 K. Kotheekar,<sup>60</sup> S. Pandey,<sup>60</sup> A. Rane,<sup>60</sup> S. Sharma,<sup>60</sup> S. Chenarani,<sup>61,bb</sup> E. Eskandari Tadavani,<sup>61</sup> S. M. Etesami,<sup>61,bb</sup>  
 M. Khakzad,<sup>61</sup> M. Mohammadi Najafabadi,<sup>61</sup> M. Naseri,<sup>61</sup> S. Paktinat Mehdiabadi,<sup>61,cc</sup> F. Rezaei Hosseinabadi,<sup>61</sup>  
 B. Safarzadeh,<sup>61,dd</sup> M. Zeinali,<sup>61</sup> M. Felcini,<sup>62</sup> M. Grunewald,<sup>62</sup> M. Abbrescia,<sup>63a,63b</sup> C. Calabria,<sup>63a,63b</sup> A. Colaleo,<sup>63a</sup>  
 D. Creanza,<sup>63a,63c</sup> L. Cristella,<sup>63a,63b</sup> N. De Filippis,<sup>63a,63c</sup> M. De Palma,<sup>63a,63b</sup> F. Errico,<sup>63a,63b</sup> L. Fiore,<sup>63a</sup> G. Iaselli,<sup>63a,63c</sup>  
 S. Lezki,<sup>63a,63b</sup> G. Maggi,<sup>63a,63c</sup> M. Maggi,<sup>63a</sup> G. Miniello,<sup>63a,63b</sup> S. My,<sup>63a,63b</sup> S. Nuzzo,<sup>63a,63b</sup> A. Pompili,<sup>63a,63b</sup>  
 G. Pugliese,<sup>63a,63c</sup> R. Radogna,<sup>63a</sup> A. Ranieri,<sup>63a</sup> G. Selvaggi,<sup>63a,63b</sup> A. Sharma,<sup>63a</sup> L. Silvestris,<sup>63a,q</sup> R. Venditti,<sup>63a</sup>  
 P. Verwilligen,<sup>63a</sup> G. Abbiendi,<sup>64a</sup> C. Battilana,<sup>64a,64b</sup> D. Bonacorsi,<sup>64a,64b</sup> L. Borgonovi,<sup>64a,64b</sup> S. Braibant-Giacomelli,<sup>64a,64b</sup>  
 R. Campanini,<sup>64a,64b</sup> P. Capiluppi,<sup>64a,64b</sup> A. Castro,<sup>64a,64b</sup> F. R. Cavallo,<sup>64a</sup> S. S. Chhibra,<sup>64a</sup> G. Codispoti,<sup>64a,64b</sup>



M. Cuffiani,<sup>64a,64b</sup> G. M. Dallavalle,<sup>64a</sup> F. Fabbri,<sup>64a</sup> A. Fanfani,<sup>64a,64b</sup> D. Fasanella,<sup>64a,64b</sup> P. Giacomelli,<sup>64a</sup> C. Grandi,<sup>64a</sup>  
 L. Guiducci,<sup>64a,64b</sup> S. Marcellini,<sup>64a</sup> G. Masetti,<sup>64a</sup> A. Montanari,<sup>64a</sup> F. L. Navarra,<sup>64a,64b</sup> A. Perrotta,<sup>64a</sup> A. M. Rossi,<sup>64a,64b</sup>  
 T. Rovelli,<sup>64a,64b</sup> G. P. Siroli,<sup>64a,64b</sup> N. Tosi,<sup>64a</sup> S. Albergo,<sup>65a,65b</sup> S. Costa,<sup>65a,65b</sup> A. Di Mattia,<sup>65a</sup> F. Giordano,<sup>65a,65b</sup>  
 R. Potenza,<sup>65a,65b</sup> A. Tricomi,<sup>65a,65b</sup> C. Tuve,<sup>65a,65b</sup> G. Barbagli,<sup>66a</sup> K. Chatterjee,<sup>66a,66b</sup> V. Ciulli,<sup>66a,66b</sup> C. Civinini,<sup>66a</sup>  
 R. D'Alessandro,<sup>66a,66b</sup> E. Focardi,<sup>66a,66b</sup> P. Lenzi,<sup>66a,66b</sup> M. Meschini,<sup>66a</sup> S. Paoletti,<sup>66a</sup> L. Russo,<sup>66a,ee</sup> G. Sguazzoni,<sup>66a</sup>  
 D. Strom,<sup>66a</sup> L. Viliani,<sup>66a,66b,q</sup> L. Benussi,<sup>67</sup> S. Bianco,<sup>67</sup> F. Fabbri,<sup>67</sup> D. Piccolo,<sup>67</sup> F. Primavera,<sup>67,q</sup> V. Calvelli,<sup>68a,68b</sup>  
 F. Ferro,<sup>68a</sup> E. Robutti,<sup>68a</sup> S. Tosi,<sup>68a,68b</sup> A. Benaglia,<sup>69a</sup> L. Brianza,<sup>69a,69b</sup> F. Brivio,<sup>69a,69b</sup> V. Ciriolo,<sup>69a,69b</sup>  
 M. E. Dinardo,<sup>69a,69b</sup> S. Fiorendi,<sup>69a,69b</sup> S. Gennai,<sup>69a</sup> A. Ghezzi,<sup>69a,69b</sup> P. Govoni,<sup>69a,69b</sup> M. Malberti,<sup>69a,69b</sup> S. Malvezzi,<sup>69a</sup>  
 R. A. Manzoni,<sup>69a,69b</sup> D. Menasce,<sup>69a</sup> L. Moroni,<sup>69a</sup> M. Paganoni,<sup>69a,69b</sup> K. Pauwels,<sup>69a,69b</sup> D. Pedrini,<sup>69a</sup> S. Pigazzini,<sup>69a,69b,ff</sup>  
 S. Ragazzi,<sup>69a,69b</sup> N. Redaelli,<sup>69a</sup> T. Tabarelli de Fatis,<sup>69a,69b</sup> S. Buontempo,<sup>70a</sup> N. Cavallo,<sup>70a,70c</sup> S. Di Guida,<sup>70a,70d,q</sup>  
 F. Fabozzi,<sup>70a,70c</sup> F. Fienga,<sup>70a,70b</sup> A. O. M. Iorio,<sup>70a,70b</sup> W. A. Khan,<sup>70a</sup> L. Lista,<sup>70a</sup> S. Meola,<sup>70a,70d,q</sup> P. Paolucci,<sup>70a,q</sup>  
 C. Sciacca,<sup>70a,70b</sup> F. Thyssen,<sup>70a</sup> P. Azzi,<sup>71a</sup> N. Bacchetta,<sup>71a</sup> L. Benato,<sup>71a,71b</sup> D. Bisello,<sup>71a,71b</sup> A. Boletti,<sup>71a,71b</sup>  
 R. Carlin,<sup>71a,71b</sup> A. Carvalho Antunes De Oliveira,<sup>71a,71b</sup> P. Checchia,<sup>71a</sup> M. Dall'Osso,<sup>71a,71b</sup> P. De Castro Manzano,<sup>71a</sup>  
 T. Dorigo,<sup>71a</sup> U. Dosselli,<sup>71a</sup> A. Gozzelino,<sup>71a</sup> S. Lacaprara,<sup>71a</sup> P. Lujan,<sup>71a</sup> M. Margoni,<sup>71a,71b</sup> A. T. Meneguzzo,<sup>71a,71b</sup>  
 F. Montecassiano,<sup>71a</sup> N. Pozzobon,<sup>71a,71b</sup> P. Ronchese,<sup>71a,71b</sup> R. Rossin,<sup>71a,71b</sup> F. Simonetto,<sup>71a,71b</sup> E. Torassa,<sup>71a</sup>  
 M. Zanetti,<sup>71a,71b</sup> P. Zotto,<sup>71a,71b</sup> G. Zumerle,<sup>71a,71b</sup> A. Braghieri,<sup>72a</sup> A. Magnani,<sup>72a</sup> P. Montagna,<sup>72a,72b</sup> S. P. Ratti,<sup>72a,72b</sup>  
 V. Re,<sup>72a</sup> M. Ressegotti,<sup>72a,72b</sup> C. Riccardi,<sup>72a,72b</sup> P. Salvini,<sup>72a</sup> I. Vai,<sup>72a,72b</sup> P. Vitulo,<sup>72a,72b</sup> L. Alunni Solestizi,<sup>73a,73b</sup>  
 M. Biasini,<sup>73a,73b</sup> G. M. Bilei,<sup>73a</sup> C. Cecchi,<sup>73a,73b</sup> D. Ciangottini,<sup>73a,73b</sup> L. Fanò,<sup>73a,73b</sup> P. Lariccia,<sup>73a,73b</sup> R. Leonardi,<sup>73a,73b</sup>  
 E. Manoni,<sup>73a</sup> G. Mantovani,<sup>73a,73b</sup> V. Mariani,<sup>73a,73b</sup> M. Menichelli,<sup>73a</sup> A. Rossi,<sup>73a,73b</sup> A. Santocchia,<sup>73a,73b</sup> D. Spiga,<sup>73a</sup>  
 K. Androsov,<sup>74a</sup> P. Azzurri,<sup>74a,q</sup> G. Bagliesi,<sup>74a</sup> T. Boccali,<sup>74a</sup> L. Borrello,<sup>74a</sup> R. Castaldi,<sup>74a</sup> M. A. Ciocci,<sup>74a,74b</sup>  
 R. Dell'Orso,<sup>74a</sup> G. Fedi,<sup>74a</sup> L. Giannini,<sup>74a,74c</sup> A. Giassi,<sup>74a</sup> M. T. Grippo,<sup>74a,ee</sup> F. Ligabue,<sup>74a,74c</sup> T. Lomtadze,<sup>74a</sup>  
 E. Manca,<sup>74a,74c</sup> G. Mandorli,<sup>74a,74c</sup> L. Martini,<sup>74a,74b</sup> A. Messineo,<sup>74a,74b</sup> F. Palla,<sup>74a</sup> A. Rizzi,<sup>74a,74b</sup> A. Savoy-Navarro,<sup>74a,gg</sup>  
 P. Spagnolo,<sup>74a</sup> R. Tenchini,<sup>74a</sup> G. Tonelli,<sup>74a,74b</sup> A. Venturi,<sup>74a</sup> P. G. Verdini,<sup>74a</sup> L. Barone,<sup>75a,75b</sup> F. Cavallari,<sup>75a</sup>  
 M. Cipriani,<sup>75a,75b</sup> N. Daci,<sup>75a</sup> D. Del Re,<sup>75a,75b,q</sup> E. Di Marco,<sup>75a,75b</sup> M. Diemoz,<sup>75a</sup> S. Gelli,<sup>75a,75b</sup> E. Longo,<sup>75a,75b</sup>  
 F. Margaroli,<sup>75a,75b</sup> B. Marzocchi,<sup>75a,75b</sup> P. Meridiani,<sup>75a</sup> G. Organtini,<sup>75a,75b</sup> R. Paramatti,<sup>75a,75b</sup> F. Preiato,<sup>75a,75b</sup>  
 S. Rahatlou,<sup>75a,75b</sup> C. Rovelli,<sup>75a</sup> F. Santanastasio,<sup>75a,75b</sup> N. Amapane,<sup>76a,76b</sup> R. Arcidiacono,<sup>76a,76c</sup> S. Argiro,<sup>76a,76b</sup>  
 M. Arneodo,<sup>76a,76c</sup> N. Bartosik,<sup>76a</sup> R. Bellan,<sup>76a,76b</sup> C. Biino,<sup>76a</sup> N. Cartiglia,<sup>76a</sup> F. Cenna,<sup>76a,76b</sup> M. Costa,<sup>76a,76b</sup>  
 R. Covarelli,<sup>76a,76b</sup> A. Degano,<sup>76a,76b</sup> N. Demaria,<sup>76a</sup> B. Kiani,<sup>76a,76b</sup> C. Mariotti,<sup>76a</sup> S. Maselli,<sup>76a</sup> E. Migliore,<sup>76a,76b</sup>  
 V. Monaco,<sup>76a,76b</sup> E. Monteil,<sup>76a,76b</sup> M. Monteno,<sup>76a</sup> M. M. Obertino,<sup>76a,76b</sup> L. Pacher,<sup>76a,76b</sup> N. Pastrone,<sup>76a</sup> M. Pelliccioni,<sup>76a</sup>  
 G. L. Pinna Angioni,<sup>76a,76b</sup> F. Ravera,<sup>76a,76b</sup> A. Romero,<sup>76a,76b</sup> M. Ruspa,<sup>76a,76c</sup> R. Sacchi,<sup>76a,76b</sup> K. Shchelina,<sup>76a,76b</sup> V. Sola,<sup>76a</sup>  
 A. Solano,<sup>76a,76b</sup> A. Staiano,<sup>76a</sup> P. Traczyk,<sup>76a,76b</sup> S. Belforte,<sup>77a</sup> M. Casarsa,<sup>77a</sup> F. Cossutti,<sup>77a</sup> G. Della Ricca,<sup>77a,77b</sup>  
 A. Zanetti,<sup>77a</sup> D. H. Kim,<sup>78</sup> G. N. Kim,<sup>78</sup> M. S. Kim,<sup>78</sup> J. Lee,<sup>78</sup> S. Lee,<sup>78</sup> S. W. Lee,<sup>78</sup> C. S. Moon,<sup>78</sup> Y. D. Oh,<sup>78</sup> S. Sekmen,<sup>78</sup>  
 D. C. Son,<sup>78</sup> Y. C. Yang,<sup>78</sup> A. Lee,<sup>79</sup> H. Kim,<sup>80</sup> D. H. Moon,<sup>80</sup> G. Oh,<sup>80</sup> J. A. Brochero Cifuentes,<sup>81</sup> J. Goh,<sup>81</sup> T. J. Kim,<sup>81</sup>  
 S. Cho,<sup>82</sup> S. Choi,<sup>82</sup> Y. Go,<sup>82</sup> D. Gyun,<sup>82</sup> S. Ha,<sup>82</sup> B. Hong,<sup>82</sup> Y. Jo,<sup>82</sup> Y. Kim,<sup>82</sup> K. Lee,<sup>82</sup> K. S. Lee,<sup>82</sup> S. Lee,<sup>82</sup> J. Lim,<sup>82</sup>  
 S. K. Park,<sup>82</sup> Y. Roh,<sup>82</sup> J. Almond,<sup>83</sup> J. Kim,<sup>83</sup> J. S. Kim,<sup>83</sup> H. Lee,<sup>83</sup> K. Lee,<sup>83</sup> K. Nam,<sup>83</sup> S. B. Oh,<sup>83</sup> B. C. Radburn-Smith,<sup>83</sup>  
 S. h. Seo,<sup>83</sup> U. K. Yang,<sup>83</sup> H. D. Yoo,<sup>83</sup> G. B. Yu,<sup>83</sup> M. Choi,<sup>84</sup> H. Kim,<sup>84</sup> J. H. Kim,<sup>84</sup> J. S. H. Lee,<sup>84</sup> I. C. Park,<sup>84</sup> Y. Choi,<sup>85</sup>  
 C. Hwang,<sup>85</sup> J. Lee,<sup>85</sup> I. Yu,<sup>85</sup> V. Dudenias,<sup>86</sup> A. Juodagalvis,<sup>86</sup> J. Vaitkus,<sup>86</sup> I. Ahmed,<sup>87</sup> Z. A. Ibrahim,<sup>87</sup>  
 M. A. B. Md Ali,<sup>87,hh</sup> F. Mohamad Idris,<sup>87,ii</sup> W. A. T. Wan Abdullah,<sup>87</sup> M. N. Yusli,<sup>87</sup> Z. Zolkapli,<sup>87</sup> R. Reyes-Almanza,<sup>88</sup>  
 G. Ramirez-Sanchez,<sup>88</sup> M. C. Duran-Osuna,<sup>88</sup> H. Castilla-Valdez,<sup>88</sup> E. De La Cruz-Burelo,<sup>88</sup> I. Heredia-De La Cruz,<sup>88,ij</sup>  
 R. I. Rabadan-Trejo,<sup>88</sup> R. Lopez-Fernandez,<sup>88</sup> J. Mejia Guisao,<sup>88</sup> A. Sanchez-Hernandez,<sup>88</sup> S. Carrillo Moreno,<sup>89</sup>  
 C. Oropeza Barrera,<sup>89</sup> F. Vazquez Valencia,<sup>89</sup> I. Pedraza,<sup>90</sup> H. A. Salazar Ibarguen,<sup>90</sup> C. Uribe Estrada,<sup>90</sup>  
 A. Morelos Pineda,<sup>91</sup> D. Krofcheck,<sup>92</sup> P. H. Butler,<sup>93</sup> A. Ahmad,<sup>94</sup> M. Ahmad,<sup>94</sup> Q. Hassan,<sup>94</sup> H. R. Hoorani,<sup>94</sup>  
 A. Saddique,<sup>94</sup> M. A. Shah,<sup>94</sup> M. Shoaib,<sup>94</sup> M. Waqas,<sup>94</sup> H. Bialkowska,<sup>95</sup> M. Bluj,<sup>95</sup> B. Boimska,<sup>95</sup> T. Frueboes,<sup>95</sup>  
 M. Górski,<sup>95</sup> M. Kazana,<sup>95</sup> K. Nawrocki,<sup>95</sup> M. Szleper,<sup>95</sup> P. Zalewski,<sup>95</sup> K. Bunkowski,<sup>96</sup> A. Byszuk,<sup>96,kk</sup> K. Doroba,<sup>96</sup>  
 A. Kalinowski,<sup>96</sup> M. Konecki,<sup>96</sup> J. Krolikowski,<sup>96</sup> M. Misiura,<sup>96</sup> M. Olszewski,<sup>96</sup> A. Pyskir,<sup>96</sup> M. Walczak,<sup>96</sup> P. Bargassa,<sup>97</sup>  
 C. Beirão Da Cruz E Silva,<sup>97</sup> A. Di Francesco,<sup>97</sup> P. Faccioli,<sup>97</sup> B. Galinhas,<sup>97</sup> M. Gallinaro,<sup>97</sup> J. Hollar,<sup>97</sup> N. Leonardo,<sup>97</sup>  
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 V. Matveev,<sup>98,ll,mmm</sup> V. Palichik,<sup>98</sup> V. Perelygin,<sup>98</sup> S. Shmatov,<sup>98</sup> S. Shulha,<sup>98</sup> N. Skatchkov,<sup>98</sup> V. Smirnov,<sup>98</sup> N. Voytishin,<sup>98</sup>

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Kovitanggoon,<sup>119</sup> G. Singh,<sup>119</sup> N. Srimanobhas,<sup>119</sup> F. Boran,<sup>120</sup> S. Cerci,<sup>120,yy</sup> S. Damarseckin,<sup>120</sup> Z. S. Demiroglu,<sup>120</sup> C. Dozen,<sup>120</sup> I. Dumanoglu,<sup>120</sup> S. Girgis,<sup>120</sup> G. Gokbulut,<sup>120</sup> Y. Guler,<sup>120</sup> I. Hos,<sup>120,zz</sup> E. E. Kangal,<sup>120,aaa</sup> O. Kara,<sup>120</sup> A. Kayis Topaksu,<sup>120</sup> U. Kiminsu,<sup>120</sup> M. Oglakci,<sup>120</sup> G. Onengut,<sup>120,bbb</sup> K. Ozdemir,<sup>120,ccc</sup> D. Sunar Cerci,<sup>120,yy</sup> B. Tali,<sup>120,yy</sup> S. Turkcapar,<sup>120</sup> I. S. Zorbakir,<sup>120</sup> C. Zorbilmez,<sup>120</sup> B. Bilin,<sup>121</sup> G. Karapinar,<sup>121,ddd</sup> K. Ocalan,<sup>121,eee</sup> M. Yalvac,<sup>121</sup> M. Zeyrek,<sup>121</sup> E. Gülmez,<sup>122</sup> M. Kaya,<sup>122,fff</sup> O. Kaya,<sup>122,ggg</sup> S. Tekten,<sup>122</sup> E. A. Yetkin,<sup>122,hhh</sup> M. N. Agaras,<sup>123</sup>

S. Atay,<sup>123</sup> A. Cakir,<sup>123</sup> K. Cankocak,<sup>123</sup> B. Gryniov,<sup>124</sup> L. Levchuk,<sup>125</sup> F. Ball,<sup>126</sup> L. Beck,<sup>126</sup> J. J. Brooke,<sup>126</sup> D. Burns,<sup>126</sup>  
 E. Clement,<sup>126</sup> D. Cussans,<sup>126</sup> O. Davignon,<sup>126</sup> H. Flacher,<sup>126</sup> J. Goldstein,<sup>126</sup> G. P. Heath,<sup>126</sup> H. F. Heath,<sup>126</sup> J. Jacob,<sup>126</sup>  
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 M. Masciovecchio,<sup>138</sup> D. Olivito,<sup>138</sup> S. Padhi,<sup>138</sup> M. Pieri,<sup>138</sup> M. Sani,<sup>138</sup> V. Sharma,<sup>138</sup> S. Simon,<sup>138</sup> M. Tadel,<sup>138</sup>  
 A. Vartak,<sup>138</sup> S. Wasserbaech,<sup>138,iii</sup> J. Wood,<sup>138</sup> F. Würthwein,<sup>138</sup> A. Yagil,<sup>138</sup> G. Zevi Della Porta,<sup>138</sup> N. Amin,<sup>139</sup>  
 R. Bhandari,<sup>139</sup> J. Bradmiller-Feld,<sup>139</sup> C. Campagnari,<sup>139</sup> A. Dishaw,<sup>139</sup> V. Dutta,<sup>139</sup> M. Franco Sevilla,<sup>139</sup> C. George,<sup>139</sup>  
 F. Golf,<sup>139</sup> L. Gouskos,<sup>139</sup> J. Gran,<sup>139</sup> R. Heller,<sup>139</sup> J. Incandela,<sup>139</sup> S. D. Mullin,<sup>139</sup> A. Ovcharova,<sup>139</sup> H. Qu,<sup>139</sup>  
 J. Richman,<sup>139</sup> D. Stuart,<sup>139</sup> I. Suarez,<sup>139</sup> J. Yoo,<sup>139</sup> D. Anderson,<sup>140</sup> J. Bendavid,<sup>140</sup> A. Bornheim,<sup>140</sup> J. M. Lawhorn,<sup>140</sup>  
 H. B. Newman,<sup>140</sup> T. Nguyen,<sup>140</sup> C. Pena,<sup>140</sup> M. Spiropulu,<sup>140</sup> J. R. Vlimant,<sup>140</sup> S. Xie,<sup>140</sup> Z. Zhang,<sup>140</sup> R. Y. Zhu,<sup>140</sup>  
 M. B. Andrews,<sup>141</sup> T. Ferguson,<sup>141</sup> T. Mudholkar,<sup>141</sup> M. Paulini,<sup>141</sup> J. Russ,<sup>141</sup> M. Sun,<sup>141</sup> H. Vogel,<sup>141</sup> I. Vorobiev,<sup>141</sup>  
 M. Weinberg,<sup>141</sup> J. P. Cumalat,<sup>142</sup> W. T. Ford,<sup>142</sup> F. Jensen,<sup>142</sup> A. Johnson,<sup>142</sup> M. Krohn,<sup>142</sup> S. Leontsinis,<sup>142</sup>  
 T. Mulholland,<sup>142</sup> K. Stenson,<sup>142</sup> S. R. Wagner,<sup>142</sup> J. Alexander,<sup>143</sup> J. Chaves,<sup>143</sup> J. Chu,<sup>143</sup> S. Dittmer,<sup>143</sup> K. McDermott,<sup>143</sup>  
 N. Mirman,<sup>143</sup> J. R. Patterson,<sup>143</sup> D. Quach,<sup>143</sup> A. Rinkevicius,<sup>143</sup> A. Ryd,<sup>143</sup> L. Skinnari,<sup>143</sup> L. Soffi,<sup>143</sup> S. M. Tan,<sup>143</sup>  
 Z. Tao,<sup>143</sup> J. Thom,<sup>143</sup> J. Tucker,<sup>143</sup> P. Wittich,<sup>143</sup> M. Zientek,<sup>143</sup> S. Abdullin,<sup>144</sup> M. Albrow,<sup>144</sup> M. Alyari,<sup>144</sup> G. Apollinari,<sup>144</sup>  
 A. Apresyan,<sup>144</sup> A. Apyan,<sup>144</sup> S. Banerjee,<sup>144</sup> L. A. T. Bauerdick,<sup>144</sup> A. Beretvas,<sup>144</sup> J. Berryhill,<sup>144</sup> P. C. Bhat,<sup>144</sup>  
 G. Bolla,<sup>144,a</sup> K. Burkett,<sup>144</sup> J. N. Butler,<sup>144</sup> A. Canepa,<sup>144</sup> G. B. Cerati,<sup>144</sup> H. W. K. Cheung,<sup>144</sup> F. Chlebana,<sup>144</sup>  
 M. Cremonesi,<sup>144</sup> J. Duarte,<sup>144</sup> V. D. Elvira,<sup>144</sup> J. Freeman,<sup>144</sup> Z. Gecse,<sup>144</sup> E. Gottschalk,<sup>144</sup> L. Gray,<sup>144</sup> D. Green,<sup>144</sup>  
 S. Grünendahl,<sup>144</sup> O. Gutsche,<sup>144</sup> R. M. Harris,<sup>144</sup> S. Hasegawa,<sup>144</sup> J. Hirschauer,<sup>144</sup> Z. Hu,<sup>144</sup> B. Jayatilaka,<sup>144</sup>  
 S. Jindariani,<sup>144</sup> M. Johnson,<sup>144</sup> U. Joshi,<sup>144</sup> B. Klima,<sup>144</sup> B. Kreis,<sup>144</sup> S. Lammel,<sup>144</sup> D. Lincoln,<sup>144</sup> R. Lipton,<sup>144</sup> M. Liu,<sup>144</sup>  
 T. Liu,<sup>144</sup> R. Lopes De Sá,<sup>144</sup> J. Lykken,<sup>144</sup> K. Maeshima,<sup>144</sup> N. Magini,<sup>144</sup> J. M. Marraffino,<sup>144</sup> D. Mason,<sup>144</sup> P. McBride,<sup>144</sup>  
 P. Merkel,<sup>144</sup> S. Mrenna,<sup>144</sup> S. Nahn,<sup>144</sup> V. O'Dell,<sup>144</sup> K. Pedro,<sup>144</sup> O. Prokofyev,<sup>144</sup> G. Rakness,<sup>144</sup> L. Ristori,<sup>144</sup>  
 B. Schneider,<sup>144</sup> E. Sexton-Kennedy,<sup>144</sup> A. Soha,<sup>144</sup> W. J. Spalding,<sup>144</sup> L. Spiegel,<sup>144</sup> S. Stoynev,<sup>144</sup> J. Strait,<sup>144</sup> N. Strobbe,<sup>144</sup>  
 L. Taylor,<sup>144</sup> S. Tkaczyk,<sup>144</sup> N. V. Tran,<sup>144</sup> L. Uplegger,<sup>144</sup> E. W. Vaandering,<sup>144</sup> C. Vernieri,<sup>144</sup> M. Verzocchi,<sup>144</sup> R. Vidal,<sup>144</sup>  
 M. Wang,<sup>144</sup> H. A. Weber,<sup>144</sup> A. Whitbeck,<sup>144</sup> D. Acosta,<sup>145</sup> P. Avery,<sup>145</sup> P. Bortignon,<sup>145</sup> D. Bourilkov,<sup>145</sup> A. Brinkerhoff,<sup>145</sup>  
 A. Carnes,<sup>145</sup> M. Carver,<sup>145</sup> D. Curry,<sup>145</sup> R. D. Field,<sup>145</sup> I. K. Furic,<sup>145</sup> S. V. Gleyzer,<sup>145</sup> B. M. Joshi,<sup>145</sup> J. Konigsberg,<sup>145</sup>  
 A. Korytov,<sup>145</sup> K. Kotov,<sup>145</sup> P. Ma,<sup>145</sup> K. Matchev,<sup>145</sup> H. Mei,<sup>145</sup> G. Mitselmakher,<sup>145</sup> D. Rank,<sup>145</sup> K. Shi,<sup>145</sup> D. Sperka,<sup>145</sup>  
 N. Terentyev,<sup>145</sup> L. Thomas,<sup>145</sup> J. Wang,<sup>145</sup> S. Wang,<sup>145</sup> J. Yelton,<sup>145</sup> Y. R. Joshi,<sup>146</sup> S. Linn,<sup>146</sup> P. Markowitz,<sup>146</sup>

J. L. Rodriguez,<sup>146</sup> A. Ackert,<sup>147</sup> T. Adams,<sup>147</sup> A. Askew,<sup>147</sup> S. Hagopian,<sup>147</sup> V. Hagopian,<sup>147</sup> K. F. Johnson,<sup>147</sup> T. Kolberg,<sup>147</sup> G. Martinez,<sup>147</sup> T. Perry,<sup>147</sup> H. Prosper,<sup>147</sup> A. Saha,<sup>147</sup> A. Santra,<sup>147</sup> V. Sharma,<sup>147</sup> R. Yohay,<sup>147</sup> M. M. Baarmand,<sup>148</sup> V. Bhopatkar,<sup>148</sup> S. Colafranceschi,<sup>148</sup> M. Hohlmann,<sup>148</sup> D. Noonan,<sup>148</sup> T. Roy,<sup>148</sup> F. Yumiceva,<sup>148</sup> M. R. Adams,<sup>149</sup> L. Apanasevich,<sup>149</sup> D. Berry,<sup>149</sup> R. R. Betts,<sup>149</sup> R. Cavanaugh,<sup>149</sup> X. Chen,<sup>149</sup> O. Evdokimov,<sup>149</sup> C. E. Gerber,<sup>149</sup> D. A. Hangal,<sup>149</sup> D. J. Hofman,<sup>149</sup> K. Jung,<sup>149</sup> J. Kamin,<sup>149</sup> I. D. Sandoval Gonzalez,<sup>149</sup> M. B. Tonjes,<sup>149</sup> H. Trauger,<sup>149</sup> N. Varelas,<sup>149</sup> H. Wang,<sup>149</sup> Z. Wu,<sup>149</sup> J. Zhang,<sup>149</sup> B. Bilki,<sup>150,mmm</sup> W. Clarida,<sup>150</sup> K. Dilsiz,<sup>150,nnn</sup> S. Durgut,<sup>150</sup> R. P. Gandrajula,<sup>150</sup> M. Haytmyradov,<sup>150</sup> V. Khristenko,<sup>150</sup> J.-P. Merlo,<sup>150</sup> H. Mermerkaya,<sup>150,ooo</sup> A. Mestvirishvili,<sup>150</sup> A. Moeller,<sup>150</sup> J. Nachtman,<sup>150</sup> H. Ogul,<sup>150,ppp</sup> Y. Onel,<sup>150</sup> F. Ozok,<sup>150,qqq</sup> A. Penzo,<sup>150</sup> C. Snyder,<sup>150</sup> E. Tiras,<sup>150</sup> J. Wetzel,<sup>150</sup> K. Yi,<sup>150</sup> B. Blumenfeld,<sup>151</sup> A. Cocoros,<sup>151</sup> N. Eminizer,<sup>151</sup> D. Fehling,<sup>151</sup> L. Feng,<sup>151</sup> A. V. Gritsan,<sup>151</sup> P. Maksimovic,<sup>151</sup> J. Roskes,<sup>151</sup> U. Sarica,<sup>151</sup> M. Swartz,<sup>151</sup> M. Xiao,<sup>151</sup> C. You,<sup>151</sup> A. Al-bataineh,<sup>152</sup> P. Baringer,<sup>152</sup> A. Bean,<sup>152</sup> S. Boren,<sup>152</sup> J. Bowen,<sup>152</sup> J. Castle,<sup>152</sup> S. Khalil,<sup>152</sup> A. Kropivnitskaya,<sup>152</sup> D. Majumder,<sup>152</sup> W. Mcbrayer,<sup>152</sup> M. Murray,<sup>152</sup> C. Royon,<sup>152</sup> S. Sanders,<sup>152</sup> E. Schmitz,<sup>152</sup> J. D. Tapia Takaki,<sup>152</sup> Q. Wang,<sup>152</sup> A. Ivanov,<sup>153</sup> K. Kaadze,<sup>153</sup> Y. Maravin,<sup>153</sup> A. Mohammadi,<sup>153</sup> L. K. Saini,<sup>153</sup> N. Skhirtladze,<sup>153</sup> S. Toda,<sup>153</sup> F. Rebassoo,<sup>154</sup> D. Wright,<sup>154</sup> C. Anelli,<sup>155</sup> A. Baden,<sup>155</sup> O. Baron,<sup>155</sup> A. Belloni,<sup>155</sup> B. Calvert,<sup>155</sup> S. C. Eno,<sup>155</sup> Y. Feng,<sup>155</sup> C. Ferraioli,<sup>155</sup> N. J. Hadley,<sup>155</sup> S. Jabeen,<sup>155</sup> G. Y. Jeng,<sup>155</sup> R. G. Kellogg,<sup>155</sup> J. Kunkle,<sup>155</sup> A. C. Mignerey,<sup>155</sup> F. Ricci-Tam,<sup>155</sup> Y. H. Shin,<sup>155</sup> A. Skuja,<sup>155</sup> S. C. Tonwar,<sup>155</sup> D. Abercrombie,<sup>156</sup> B. Allen,<sup>156</sup> V. Azzolini,<sup>156</sup> R. Barbieri,<sup>156</sup> A. Baty,<sup>156</sup> R. Bi,<sup>156</sup> S. Brandt,<sup>156</sup> W. Busza,<sup>156</sup> I. A. Cali,<sup>156</sup> M. D'Alfonso,<sup>156</sup> Z. Demiragli,<sup>156</sup> G. Gomez Ceballos,<sup>156</sup> M. Goncharov,<sup>156</sup> D. Hsu,<sup>156</sup> M. Hu,<sup>156</sup> Y. Iiyama,<sup>156</sup> G. M. Innocenti,<sup>156</sup> M. Klute,<sup>156</sup> D. Kovalskyi,<sup>156</sup> Y. S. Lai,<sup>156</sup> Y.-J. Lee,<sup>156</sup> A. Levin,<sup>156</sup> P. D. Luckey,<sup>156</sup> B. Maier,<sup>156</sup> A. C. Marini,<sup>156</sup> C. McGinn,<sup>156</sup> C. Mironov,<sup>156</sup> S. Narayanan,<sup>156</sup> X. Niu,<sup>156</sup> C. Paus,<sup>156</sup> C. Roland,<sup>156</sup> G. Roland,<sup>156</sup> J. Salfeld-Nebgen,<sup>156</sup> G. S. F. Stephans,<sup>156</sup> K. Tatar,<sup>156</sup> D. Velicanu,<sup>156</sup> J. Wang,<sup>156</sup> T. W. Wang,<sup>156</sup> B. Wyslouch,<sup>156</sup> A. C. Benvenuti,<sup>157</sup> R. M. Chatterjee,<sup>157</sup> A. Evans,<sup>157</sup> P. Hansen,<sup>157</sup> J. Hiltbrand,<sup>157</sup> S. Kalafut,<sup>157</sup> Y. Kubota,<sup>157</sup> Z. Lesko,<sup>157</sup> J. Mans,<sup>157</sup> S. Nourbakhsh,<sup>157</sup> N. Ruckstuhl,<sup>157</sup> R. Rusack,<sup>157</sup> J. Turkewitz,<sup>157</sup> M. A. Wadud,<sup>157</sup> J. G. Acosta,<sup>158</sup> S. Oliveros,<sup>158</sup> E. Avdeeva,<sup>159</sup> K. Bloom,<sup>159</sup> D. R. Claes,<sup>159</sup> C. Fangmeier,<sup>159</sup> R. Gonzalez Suarez,<sup>159</sup> R. Kamalieddin,<sup>159</sup> I. Kravchenko,<sup>159</sup> J. Monroy,<sup>159</sup> J. E. Siado,<sup>159</sup> G. R. Snow,<sup>159</sup> B. Stieger,<sup>159</sup> J. Dolen,<sup>160</sup> A. Godshalk,<sup>160</sup> C. Harrington,<sup>160</sup> I. Iashvili,<sup>160</sup> D. Nguyen,<sup>160</sup> A. Parker,<sup>160</sup> S. Rappoccio,<sup>160</sup> B. Roozbahani,<sup>160</sup> G. Alverson,<sup>161</sup> E. Barberis,<sup>161</sup> A. Hortiangtham,<sup>161</sup> A. Massironi,<sup>161</sup> D. M. Morse,<sup>161</sup> T. Orimoto,<sup>161</sup> R. Teixeira De Lima,<sup>161</sup> D. Trocino,<sup>161</sup> D. Wood,<sup>161</sup> S. Bhattacharya,<sup>162</sup> O. Charaf,<sup>162</sup> K. A. Hahn,<sup>162</sup> N. Mucia,<sup>162</sup> N. Odell,<sup>162</sup> B. Pollack,<sup>162</sup> M. H. Schmitt,<sup>162</sup> K. Sung,<sup>162</sup> M. Trovato,<sup>162</sup> M. Velasco,<sup>162</sup> N. Dev,<sup>163</sup> M. Hildreth,<sup>163</sup> K. Hurtado Anampa,<sup>163</sup> C. Jessop,<sup>163</sup> D. J. Karmgard,<sup>163</sup> N. Kellams,<sup>163</sup> K. Lannon,<sup>163</sup> N. Loukas,<sup>163</sup> N. Marinelli,<sup>163</sup> F. Meng,<sup>163</sup> C. Mueller,<sup>163</sup> Y. Musienko,<sup>163,ll</sup> M. Planer,<sup>163</sup> A. Reinsvold,<sup>163</sup> R. Ruchti,<sup>163</sup> G. Smith,<sup>163</sup> S. Taroni,<sup>163</sup> M. Wayne,<sup>163</sup> M. Wolf,<sup>163</sup> A. Woodard,<sup>163</sup> J. Alimena,<sup>164</sup> L. Antonelli,<sup>164</sup> B. Bylsma,<sup>164</sup> L. S. Durkin,<sup>164</sup> S. Flowers,<sup>164</sup> B. Francis,<sup>164</sup> A. Hart,<sup>164</sup> C. Hill,<sup>164</sup> W. Ji,<sup>164</sup> B. Liu,<sup>164</sup> W. Luo,<sup>164</sup> D. Puigh,<sup>164</sup> B. L. Winer,<sup>164</sup> H. W. Wulsin,<sup>164</sup> S. Cooperstein,<sup>165</sup> O. Driga,<sup>165</sup> P. Elmer,<sup>165</sup> J. Hardenbrook,<sup>165</sup> P. Hebda,<sup>165</sup> S. Higginbotham,<sup>165</sup> D. Lange,<sup>165</sup> J. Luo,<sup>165</sup> D. Marlow,<sup>165</sup> K. Mei,<sup>165</sup> I. Ojalvo,<sup>165</sup> J. Olsen,<sup>165</sup> C. Palmer,<sup>165</sup> P. Piroué,<sup>165</sup> D. Stickland,<sup>165</sup> C. Tully,<sup>165</sup> S. Malik,<sup>166</sup> S. Norberg,<sup>166</sup> A. Barker,<sup>167</sup> V. E. Barnes,<sup>167</sup> S. Das,<sup>167</sup> S. Folgueras,<sup>167</sup> L. Gutay,<sup>167</sup> M. K. Jha,<sup>167</sup> M. Jones,<sup>167</sup> A. W. Jung,<sup>167</sup> A. Khatiwada,<sup>167</sup> D. H. Miller,<sup>167</sup> N. Neumeister,<sup>167</sup> C. C. Peng,<sup>167</sup> H. Qiu,<sup>167</sup> J. F. Schulte,<sup>167</sup> J. Sun,<sup>167</sup> F. Wang,<sup>167</sup> W. Xie,<sup>167</sup> T. Cheng,<sup>168</sup> N. Parashar,<sup>168</sup> J. Stupak,<sup>168</sup> A. Adair,<sup>169</sup> Z. Chen,<sup>169</sup> K. M. Ecklund,<sup>169</sup> S. Freed,<sup>169</sup> F. J. M. Geurts,<sup>169</sup> M. Guilbaud,<sup>169</sup> M. Kilpatrick,<sup>169</sup> W. Li,<sup>169</sup> B. Michlin,<sup>169</sup> M. Northup,<sup>169</sup> B. P. Padley,<sup>169</sup> J. Roberts,<sup>169</sup> J. Rorie,<sup>169</sup> W. Shi,<sup>169</sup> Z. Tu,<sup>169</sup> J. Zabel,<sup>169</sup> A. Zhang,<sup>169</sup> A. Bodek,<sup>170</sup> P. de Barbaro,<sup>170</sup> R. Demina,<sup>170</sup> Y. t. Duh,<sup>170</sup> T. Ferbel,<sup>170</sup> M. Galanti,<sup>170</sup> A. Garcia-Bellido,<sup>170</sup> J. Han,<sup>170</sup> O. Hindrichs,<sup>170</sup> A. Khukhunaishvili,<sup>170</sup> K. H. Lo,<sup>170</sup> P. Tan,<sup>170</sup> M. Verzetti,<sup>170</sup> R. Ciesielski,<sup>171</sup> K. Goulianos,<sup>171</sup> C. Mesropian,<sup>171</sup> A. Agapitos,<sup>172</sup> J. P. Chou,<sup>172</sup> Y. Gershtein,<sup>172</sup> T. A. Gómez Espinosa,<sup>172</sup> E. Halkiadakis,<sup>172</sup> M. Heindl,<sup>172</sup> E. Hughes,<sup>172</sup> S. Kaplan,<sup>172</sup> R. Kunnawalkam Elayavalli,<sup>172</sup> S. Kyriacou,<sup>172</sup> A. Lath,<sup>172</sup> R. Montalvo,<sup>172</sup> K. Nash,<sup>172</sup> M. Osherson,<sup>172</sup> H. Saka,<sup>172</sup> S. Salur,<sup>172</sup> S. Schnetzer,<sup>172</sup> D. Sheffield,<sup>172</sup> S. Somalwar,<sup>172</sup> R. Stone,<sup>172</sup> S. Thomas,<sup>172</sup> P. Thomassen,<sup>172</sup> M. Walker,<sup>172</sup> A. G. Delannoy,<sup>173</sup> M. Foerster,<sup>173</sup> J. Heideman,<sup>173</sup> G. Riley,<sup>173</sup> K. Rose,<sup>173</sup> S. Spanier,<sup>173</sup> K. Thapa,<sup>173</sup> O. Bouhali,<sup>174,rrr</sup> A. Castaneda Hernandez,<sup>174,rrr</sup> A. Celik,<sup>174</sup> M. Dalchenko,<sup>174</sup> M. De Mattia,<sup>174</sup> A. Delgado,<sup>174</sup> S. Dildick,<sup>174</sup> R. Eusebi,<sup>174</sup> J. Gilmore,<sup>174</sup> T. Huang,<sup>174</sup> T. Kamon,<sup>174,sss</sup> R. Mueller,<sup>174</sup> Y. Pakhotin,<sup>174</sup> R. Patel,<sup>174</sup> A. Perloff,<sup>174</sup> L. Perniè,<sup>174</sup> D. Rathjens,<sup>174</sup> A. Safonov,<sup>174</sup> A. Tatarinov,<sup>174</sup> K. A. Ulmer,<sup>174</sup> N. Akchurin,<sup>175</sup> J. Damgov,<sup>175</sup> F. De Guio,<sup>175</sup> P. R. Duderø,<sup>175</sup> J. Faulkner,<sup>175</sup> E. Gurpinar,<sup>175</sup> S. Kunori,<sup>175</sup> K. Lamichhane,<sup>175</sup> S. W. Lee,<sup>175</sup> T. Libeiro,<sup>175</sup> T. Mengke,<sup>175</sup> S. Muthumuni,<sup>175</sup> T. Peltola,<sup>175</sup> S. Undleeb,<sup>175</sup> I. Volobouev,<sup>175</sup> Z. Wang,<sup>175</sup> S. Greene,<sup>176</sup> A. Gurrola,<sup>176</sup>

R. Janjam,<sup>176</sup> W. Johns,<sup>176</sup> C. Maguire,<sup>176</sup> A. Melo,<sup>176</sup> H. Ni,<sup>176</sup> K. Padeken,<sup>176</sup> P. Sheldon,<sup>176</sup> S. Tuo,<sup>176</sup> J. Velkovska,<sup>176</sup> Q. Xu,<sup>176</sup> M. W. Arenton,<sup>177</sup> P. Barria,<sup>177</sup> B. Cox,<sup>177</sup> R. Hirosky,<sup>177</sup> M. Joyce,<sup>177</sup> A. Ledovskoy,<sup>177</sup> H. Li,<sup>177</sup> C. Neu,<sup>177</sup> T. Sinthuprasith,<sup>177</sup> Y. Wang,<sup>177</sup> E. Wolfe,<sup>177</sup> F. Xia,<sup>177</sup> R. Harr,<sup>178</sup> P. E. Karchin,<sup>178</sup> N. Poudyal,<sup>178</sup> J. Sturdy,<sup>178</sup> P. Thapa,<sup>178</sup> S. Zaleski,<sup>178</sup> M. Brodski,<sup>179</sup> J. Buchanan,<sup>179</sup> C. Caillol,<sup>179</sup> S. Dasu,<sup>179</sup> L. Dodd,<sup>179</sup> S. Duric,<sup>179</sup> B. Gomber,<sup>179</sup> M. Grothe,<sup>179</sup> M. Herndon,<sup>179</sup> A. Hervé,<sup>179</sup> U. Hussain,<sup>179</sup> P. Klabbbers,<sup>179</sup> A. Lanaro,<sup>179</sup> A. Levine,<sup>179</sup> K. Long,<sup>179</sup> R. Loveless,<sup>179</sup> G. Polese,<sup>179</sup> T. Ruggles,<sup>179</sup> A. Savin,<sup>179</sup> N. Smith,<sup>179</sup> W. H. Smith,<sup>179</sup> D. Taylor,<sup>179</sup> and N. Woods<sup>179</sup>

(CMS Collaboration)

<sup>1</sup>*Yerevan Physics Institute, Yerevan, Armenia*

<sup>2</sup>*Institut für Hochenergiephysik, Wien, Austria*

<sup>3</sup>*Institute for Nuclear Problems, Minsk, Belarus*

<sup>4</sup>*Universiteit Antwerpen, Antwerpen, Belgium*

<sup>5</sup>*Vrije Universiteit Brussel, Brussel, Belgium*

<sup>6</sup>*Université Libre de Bruxelles, Bruxelles, Belgium*

<sup>7</sup>*Ghent University, Ghent, Belgium*

<sup>8</sup>*Université Catholique de Louvain, Louvain-la-Neuve, Belgium*

<sup>9</sup>*Université de Mons, Mons, Belgium*

<sup>10</sup>*Centro Brasileiro de Pesquisas Físicas, Rio de Janeiro, Brazil*

<sup>11</sup>*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*

<sup>12a</sup>*Universidade Estadual Paulista, São Paulo, Brazil*

<sup>12b</sup>*Universidade Federal do ABC, São Paulo, Brazil*

<sup>13</sup>*Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria*

<sup>14</sup>*University of Sofia, Sofia, Bulgaria*

<sup>15</sup>*Beihang University, Beijing, China*

<sup>16</sup>*Institute of High Energy Physics, Beijing, China*

<sup>17</sup>*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*

<sup>18</sup>*Universidad de Los Andes, Bogota, Colombia*

<sup>19</sup>*University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia*

<sup>20</sup>*University of Split, Faculty of Science, Split, Croatia*

<sup>21</sup>*Institute Rudjer Boskovic, Zagreb, Croatia*

<sup>22</sup>*University of Cyprus, Nicosia, Cyprus*

<sup>23</sup>*Charles University, Prague, Czech Republic*

<sup>24</sup>*Universidad San Francisco de Quito, Quito, Ecuador*

<sup>25</sup>*Academy of Scientific Research and Technology of the Arab Republic of Egypt,*

*Egyptian Network of High Energy Physics, Cairo, Egypt*

<sup>26</sup>*National Institute of Chemical Physics and Biophysics, Tallinn, Estonia*

<sup>27</sup>*Department of Physics, University of Helsinki, Helsinki, Finland*

<sup>28</sup>*Helsinki Institute of Physics, Helsinki, Finland*

<sup>29</sup>*Lappeenranta University of Technology, Lappeenranta, Finland*

<sup>30</sup>*IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France*

<sup>31</sup>*Laboratoire Leprince-Ringuet, Ecole polytechnique, CNRS/IN2P3, Université Paris-Saclay, Palaiseau, France*

<sup>32</sup>*Université de Strasbourg, CNRS, IPHC UMR 7178, F-67000 Strasbourg, France*

<sup>33</sup>*Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules, CNRS/IN2P3, Villeurbanne, France*

<sup>34</sup>*Université de Lyon, Université Claude Bernard Lyon 1, CNRS-IN2P3, Institut de Physique Nucléaire de Lyon, Villeurbanne, France*

<sup>35</sup>*Georgian Technical University, Tbilisi, Georgia*

<sup>36</sup>*Tbilisi State University, Tbilisi, Georgia*

<sup>37</sup>*RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany*

<sup>38</sup>*RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany*

<sup>39</sup>*RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany*

<sup>40</sup>*Deutsches Elektronen-Synchrotron, Hamburg, Germany*

<sup>41</sup>*University of Hamburg, Hamburg, Germany*

<sup>42</sup>*Institut für Experimentelle Kernphysik, Karlsruhe, Germany*

<sup>43</sup>*Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece*

<sup>44</sup>*National and Kapodistrian University of Athens, Athens, Greece*

<sup>45</sup>*National Technical University of Athens, Athens, Greece*

<sup>46</sup>*University of Ioánnina, Ioánnina, Greece*

<sup>47</sup>*MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary*

- <sup>48</sup>*Wigner Research Centre for Physics, Budapest, Hungary*  
<sup>49</sup>*Institute of Nuclear Research ATOMKI, Debrecen, Hungary*  
<sup>50</sup>*Institute of Physics, University of Debrecen, Debrecen, Hungary*  
<sup>51</sup>*Indian Institute of Science (IISc), Bangalore, India*  
<sup>52</sup>*National Institute of Science Education and Research, Bhubaneswar, India*  
<sup>53</sup>*Panjab University, Chandigarh, India*  
<sup>54</sup>*University of Delhi, Delhi, India*  
<sup>55</sup>*Saha Institute of Nuclear Physics, HBNI, Kolkata, India*  
<sup>56</sup>*Indian Institute of Technology Madras, Madras, India*  
<sup>57</sup>*Bhabha Atomic Research Centre, Mumbai, India*  
<sup>58</sup>*Tata Institute of Fundamental Research-A, Mumbai, India*  
<sup>59</sup>*Tata Institute of Fundamental Research-B, Mumbai, India*  
<sup>60</sup>*Indian Institute of Science Education and Research (IISER), Pune, India*  
<sup>61</sup>*Institute for Research in Fundamental Sciences (IPM), Tehran, Iran*  
<sup>62</sup>*University College Dublin, Dublin, Ireland*  
<sup>63a</sup>*INFN Sezione di Bari, Bari, Italy*  
<sup>63b</sup>*Università di Bari, Bari, Italy*  
<sup>63c</sup>*Politecnico di Bari, Bari, Italy*  
<sup>64a</sup>*INFN Sezione di Bologna, Bologna, Italy*  
<sup>64b</sup>*Università di Bologna, Bologna, Italy*  
<sup>65a</sup>*INFN Sezione di Catania, Catania, Italy*  
<sup>65b</sup>*Università di Catania, Catania, Italy*  
<sup>66a</sup>*INFN Sezione di Firenze, Firenze, Italy*  
<sup>66b</sup>*Università di Firenze, Firenze, Italy*  
<sup>67</sup>*INFN Laboratori Nazionali di Frascati, Frascati, Italy*  
<sup>68a</sup>*INFN Sezione di Genova, Genova, Italy*  
<sup>68b</sup>*Università di Genova, Genova, Italy*  
<sup>69a</sup>*INFN Sezione di Milano-Bicocca, Milano, Italy*  
<sup>69b</sup>*Università di Milano-Bicocca, Milano, Italy*  
<sup>70a</sup>*INFN Sezione di Napoli, Napoli, Italy*  
<sup>70b</sup>*Università di Napoli 'Federico II', Napoli, Italy*  
<sup>70c</sup>*Università della Basilicata, Potenza, Italy*  
<sup>70d</sup>*Università G. Marconi, Roma, Italy*  
<sup>71a</sup>*INFN Sezione di Padova, Padova, Italy*  
<sup>71b</sup>*Università di Padova, Padova, Italy*  
<sup>71c</sup>*Università di Trento, Trento, Italy*  
<sup>72a</sup>*INFN Sezione di Pavia, Pavia, Italy*  
<sup>72b</sup>*Università di Pavia, Pavia, Italy*  
<sup>73a</sup>*INFN Sezione di Perugia, Perugia, Italy*  
<sup>73b</sup>*Università di Perugia, Perugia, Italy*  
<sup>74a</sup>*INFN Sezione di Pisa, Pisa, Italy*  
<sup>74b</sup>*Università di Pisa, Pisa, Italy*  
<sup>74c</sup>*Scuola Normale Superiore di Pisa, Pisa, Italy*  
<sup>75a</sup>*INFN Sezione di Roma, Rome, Italy*  
<sup>75b</sup>*Sapienza Università di Roma, Rome, Italy*  
<sup>76a</sup>*INFN Sezione di Torino, Torino, Italy*  
<sup>76b</sup>*Università di Torino, Torino, Italy*  
<sup>76c</sup>*Università del Piemonte Orientale, Novara, Italy*  
<sup>77a</sup>*INFN Sezione di Trieste, Trieste, Italy*  
<sup>77b</sup>*Università di Trieste, Trieste, Italy*  
<sup>78</sup>*Kyungpook National University, Daegu, Korea*  
<sup>79</sup>*Chonbuk National University, Jeonju, Korea*  
<sup>80</sup>*Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea*  
<sup>81</sup>*Hanyang University, Seoul, Korea*  
<sup>82</sup>*Korea University, Seoul, Korea*  
<sup>83</sup>*Seoul National University, Seoul, Korea*  
<sup>84</sup>*University of Seoul, Seoul, Korea*  
<sup>85</sup>*Sungkyunkwan University, Suwon, Korea*  
<sup>86</sup>*Vilnius University, Vilnius, Lithuania*  
<sup>87</sup>*National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia*

- <sup>88</sup>*Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico*  
<sup>89</sup>*Universidad Iberoamericana, Mexico City, Mexico*  
<sup>90</sup>*Benemerita Universidad Autonoma de Puebla, Puebla, Mexico*  
<sup>91</sup>*Universidad Autónoma de San Luis Potosí, San Luis Potosí, Mexico*  
<sup>92</sup>*University of Auckland, Auckland, New Zealand*  
<sup>93</sup>*University of Canterbury, Christchurch, New Zealand*  
<sup>94</sup>*National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan*  
<sup>95</sup>*National Centre for Nuclear Research, Swierk, Poland*  
<sup>96</sup>*Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland*  
<sup>97</sup>*Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal*  
<sup>98</sup>*Joint Institute for Nuclear Research, Dubna, Russia*  
<sup>99</sup>*Petersburg Nuclear Physics Institute, Gatchina (St. Petersburg), Russia*  
<sup>100</sup>*Institute for Nuclear Research, Moscow, Russia*  
<sup>101</sup>*Institute for Theoretical and Experimental Physics, Moscow, Russia*  
<sup>102</sup>*Moscow Institute of Physics and Technology, Moscow, Russia*  
<sup>103</sup>*National Research Nuclear University 'Moscow Engineering Physics Institute' (MEPhI), Moscow, Russia*  
<sup>104</sup>*P.N. Lebedev Physical Institute, Moscow, Russia*  
<sup>105</sup>*Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia*  
<sup>106</sup>*Novosibirsk State University (NSU), Novosibirsk, Russia*  
<sup>107</sup>*State Research Center of Russian Federation, Institute for High Energy Physics, Protvino, Russia*  
<sup>108</sup>*University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia*  
<sup>109</sup>*Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain*  
<sup>110</sup>*Universidad Autónoma de Madrid, Madrid, Spain*  
<sup>111</sup>*Universidad de Oviedo, Oviedo, Spain*  
<sup>112</sup>*Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain*  
<sup>113</sup>*CERN, European Organization for Nuclear Research, Geneva, Switzerland*  
<sup>114</sup>*Paul Scherrer Institut, Villigen, Switzerland*  
<sup>115</sup>*Institute for Particle Physics, ETH Zurich, Zurich, Switzerland*  
<sup>116</sup>*Universität Zürich, Zurich, Switzerland*  
<sup>117</sup>*National Central University, Chung-Li, Taiwan*  
<sup>118</sup>*National Taiwan University (NTU), Taipei, Taiwan*  
<sup>119</sup>*Chulalongkorn University, Faculty of Science, Department of Physics, Bangkok, Thailand*  
<sup>120</sup>*Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey*  
<sup>121</sup>*Middle East Technical University, Physics Department, Ankara, Turkey*  
<sup>122</sup>*Bogazici University, Istanbul, Turkey*  
<sup>123</sup>*Istanbul Technical University, Istanbul, Turkey*  
<sup>124</sup>*Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkov, Ukraine*  
<sup>125</sup>*National Scientific Center, Kharkov Institute of Physics and Technology, Kharkov, Ukraine*  
<sup>126</sup>*University of Bristol, Bristol, United Kingdom*  
<sup>127</sup>*Rutherford Appleton Laboratory, Didcot, United Kingdom*  
<sup>128</sup>*Imperial College, London, United Kingdom*  
<sup>129</sup>*Brunel University, Uxbridge, United Kingdom*  
<sup>130</sup>*Baylor University, Waco, Texas, USA*  
<sup>131</sup>*Catholic University of America, Washington DC, USA*  
<sup>132</sup>*The University of Alabama, Tuscaloosa, Alabama, USA*  
<sup>133</sup>*Boston University, Boston, Massachusetts, USA*  
<sup>134</sup>*Brown University, Providence, Rhode Island, USA*  
<sup>135</sup>*University of California, Davis, Davis, California, USA*  
<sup>136</sup>*University of California, Los Angeles, California, USA*  
<sup>137</sup>*University of California, Riverside, Riverside, California, USA*  
<sup>138</sup>*University of California, San Diego, La Jolla, California, USA*  
<sup>139</sup>*University of California, Santa Barbara—Department of Physics, Santa Barbara, California, USA*  
<sup>140</sup>*California Institute of Technology, Pasadena, California, USA*  
<sup>141</sup>*Carnegie Mellon University, Pittsburgh, Pennsylvania, USA*  
<sup>142</sup>*University of Colorado Boulder, Boulder, Colorado, USA*  
<sup>143</sup>*Cornell University, Ithaca, New York, USA*  
<sup>144</sup>*Fermi National Accelerator Laboratory, Batavia, Illinois, USA*  
<sup>145</sup>*University of Florida, Gainesville, Florida, USA*  
<sup>146</sup>*Florida International University, Miami, Florida, USA*  
<sup>147</sup>*Florida State University, Tallahassee, Florida, USA*

- <sup>148</sup>*Florida Institute of Technology, Melbourne, Florida, USA*  
<sup>149</sup>*University of Illinois at Chicago (UIC), Chicago, Illinois, USA*  
<sup>150</sup>*The University of Iowa, Iowa City, Iowa, USA*  
<sup>151</sup>*Johns Hopkins University, Baltimore, Maryland, USA*  
<sup>152</sup>*The University of Kansas, Lawrence, Kansas, USA*  
<sup>153</sup>*Kansas State University, Manhattan, Kansas, USA*  
<sup>154</sup>*Lawrence Livermore National Laboratory, Livermore, California, USA*  
<sup>155</sup>*University of Maryland, College Park, Maryland, USA*  
<sup>156</sup>*Massachusetts Institute of Technology, Cambridge, Massachusetts, USA*  
<sup>157</sup>*University of Minnesota, Minneapolis, Minnesota, USA*  
<sup>158</sup>*University of Mississippi, Oxford, Mississippi, USA*  
<sup>159</sup>*University of Nebraska-Lincoln, Lincoln, Nebraska, USA*  
<sup>160</sup>*State University of New York at Buffalo, Buffalo, New York, USA*  
<sup>161</sup>*Northeastern University, Boston, Massachusetts, USA*  
<sup>162</sup>*Northwestern University, Evanston, Illinois, USA*  
<sup>163</sup>*University of Notre Dame, Notre Dame, Indiana, USA*  
<sup>164</sup>*The Ohio State University, Columbus, Ohio, USA*  
<sup>165</sup>*Princeton University, Princeton, New Jersey, USA*  
<sup>166</sup>*University of Puerto Rico, Mayaguez, Puerto Rico, USA*  
<sup>167</sup>*Purdue University, West Lafayette, Indiana, USA*  
<sup>168</sup>*Purdue University Northwest, Hammond, USA*  
<sup>169</sup>*Rice University, Houston, Texas, USA*  
<sup>170</sup>*University of Rochester, Rochester, New York, USA*  
<sup>171</sup>*The Rockefeller University, New York, New York, USA*  
<sup>172</sup>*Rutgers, The State University of New Jersey, Piscataway, New Jersey, USA*  
<sup>173</sup>*University of Tennessee, Knoxville, Tennessee, USA*  
<sup>174</sup>*Texas A&M University, College Station, Texas, USA*  
<sup>175</sup>*Texas Tech University, Lubbock, Texas, USA*  
<sup>176</sup>*Vanderbilt University, Nashville, Tennessee, USA*  
<sup>177</sup>*University of Virginia, Charlottesville, Virginia, USA*  
<sup>178</sup>*Wayne State University, Detroit, Michigan, USA*  
<sup>179</sup>*University of Wisconsin—Madison, Madison, Wisconsin, USA*

<sup>a</sup>Deceased.

<sup>b</sup>Also at Vienna University of Technology, Vienna, Austria.

<sup>c</sup>Also at State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China.

<sup>d</sup>Also at IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France.

<sup>e</sup>Also at Universidade Estadual de Campinas, Campinas, Brazil.

<sup>f</sup>Also at Universidade Federal de Pelotas, Pelotas, Brazil.

<sup>g</sup>Also at Université Libre de Bruxelles, Bruxelles, Belgium.

<sup>h</sup>Also at Institute for Theoretical and Experimental Physics, Moscow, Russia.

<sup>i</sup>Also at Joint Institute for Nuclear Research, Dubna, Russia.

<sup>j</sup>Also at Helwan University, Cairo, Egypt.

<sup>k</sup>Also at Zewail City of Science and Technology, Zewail, Egypt.

<sup>l</sup>Also at Fayoum University, El-Fayoum, Egypt.

<sup>m</sup>Also at British University in Egypt, Cairo, Egypt.

<sup>n</sup>Also at Ain Shams University, Cairo, Egypt.

<sup>o</sup>Also at Université de Haute Alsace, Mulhouse, France.

<sup>p</sup>Also at Skobeltsyn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia.

<sup>q</sup>Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.

<sup>r</sup>Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany.

<sup>s</sup>Also at University of Hamburg, Hamburg, Germany.

<sup>t</sup>Also at Brandenburg University of Technology, Cottbus, Germany.

<sup>u</sup>Also at MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary.

<sup>v</sup>Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.

<sup>w</sup>Also at Institute of Physics, University of Debrecen, Debrecen, Hungary.

<sup>x</sup>Also at IIT Bhubaneswar, Bhubaneswar, India.

<sup>y</sup>Also at Institute of Physics, Bhubaneswar, India.

<sup>z</sup>Also at University of Visva-Bharati, Santiniketan, India.

<sup>aa</sup>Also at University of Ruhuna, Matara, Sri Lanka.



- <sup>bb</sup> Also at Isfahan University of Technology, Isfahan, Iran.
- <sup>cc</sup> Also at Yazd University, Yazd, Iran.
- <sup>dd</sup> Also at Plasma Physics Research Center, Science and Research Branch, Islamic Azad University, Tehran, Iran.
- <sup>ee</sup> Also at Università degli Studi di Siena, Siena, Italy.
- <sup>ff</sup> Also at INFN Sezione di Milano-Bicocca, Università di Milano-Bicocca, Milano, Italy.
- <sup>gg</sup> Also at Purdue University, West Lafayette, USA.
- <sup>hh</sup> Also at International Islamic University of Malaysia, Kuala Lumpur, Malaysia.
- <sup>ii</sup> Also at Malaysian Nuclear Agency, MOSTI, Kajang, Malaysia.
- <sup>jj</sup> Also at Consejo Nacional de Ciencia y Tecnología, Mexico city, Mexico.
- <sup>kk</sup> Also at Warsaw University of Technology, Institute of Electronic Systems, Warsaw, Poland.
- <sup>ll</sup> Also at Institute for Nuclear Research, Moscow, Russia.
- <sup>mm</sup> Also at National Research Nuclear University 'Moscow Engineering Physics Institute' (MEPhI), Moscow, Russia.
- <sup>nn</sup> Also at St. Petersburg State Polytechnical University, St. Petersburg, Russia.
- <sup>oo</sup> Also at University of Florida, Gainesville, USA.
- <sup>pp</sup> Also at P.N. Lebedev Physical Institute, Moscow, Russia.
- <sup>qq</sup> Also at Budker Institute of Nuclear Physics, Novosibirsk, Russia.
- <sup>rr</sup> Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia.
- <sup>ss</sup> Also at University of Belgrade, Faculty of Physics and Vinca Institute of Nuclear Sciences, Belgrade, Serbia.
- <sup>tt</sup> Also at Scuola Normale e Sezione dell'INFN, Pisa, Italy.
- <sup>uu</sup> Also at National and Kapodistrian University of Athens, Athens, Greece.
- <sup>vv</sup> Also at Riga Technical University, Riga, Latvia.
- <sup>ww</sup> Also at Universität Zürich, Zurich, Switzerland.
- <sup>xx</sup> Also at Stefan Meyer Institute for Subatomic Physics.
- <sup>yy</sup> Also at Adiyaman University, Adiyaman, Turkey.
- <sup>zz</sup> Also at Istanbul Aydin University, Istanbul, Turkey.
- <sup>aaa</sup> Also at Mersin University, Mersin, Turkey.
- <sup>bbb</sup> Also at Cag University, Mersin, Turkey.
- <sup>ccc</sup> Also at Piri Reis University, Istanbul, Turkey.
- <sup>ddd</sup> Also at Izmir Institute of Technology, Izmir, Turkey.
- <sup>eee</sup> Also at Necmettin Erbakan University, Konya, Turkey.
- <sup>fff</sup> Also at Marmara University, Istanbul, Turkey.
- <sup>ggg</sup> Also at Kafkas University, Kars, Turkey.
- <sup>hhh</sup> Also at Istanbul Bilgi University, Istanbul, Turkey.
- <sup>iii</sup> Also at Rutherford Appleton Laboratory, Didcot, United Kingdom.
- <sup>jjj</sup> Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
- <sup>kkk</sup> Also at Instituto de Astrofísica de Canarias, La Laguna, Spain.
- <sup>lll</sup> Also at Utah Valley University, Orem, USA.
- <sup>mmm</sup> Also at Beykent University.
- <sup>nnn</sup> Also at Bingol University, Bingol, Turkey.
- <sup>ooo</sup> Also at Erzincan University, Erzincan, Turkey.
- <sup>ppp</sup> Also at Sinop University, Sinop, Turkey.
- <sup>qqq</sup> Also at Mimar Sinan University, Istanbul, Istanbul, Turkey.
- <sup>rrr</sup> Also at Texas A&M University at Qatar, Doha, Qatar.
- <sup>sss</sup> Also at Kyungpook National University, Daegu, Korea.